



Nuclear quantum optics at XFELs

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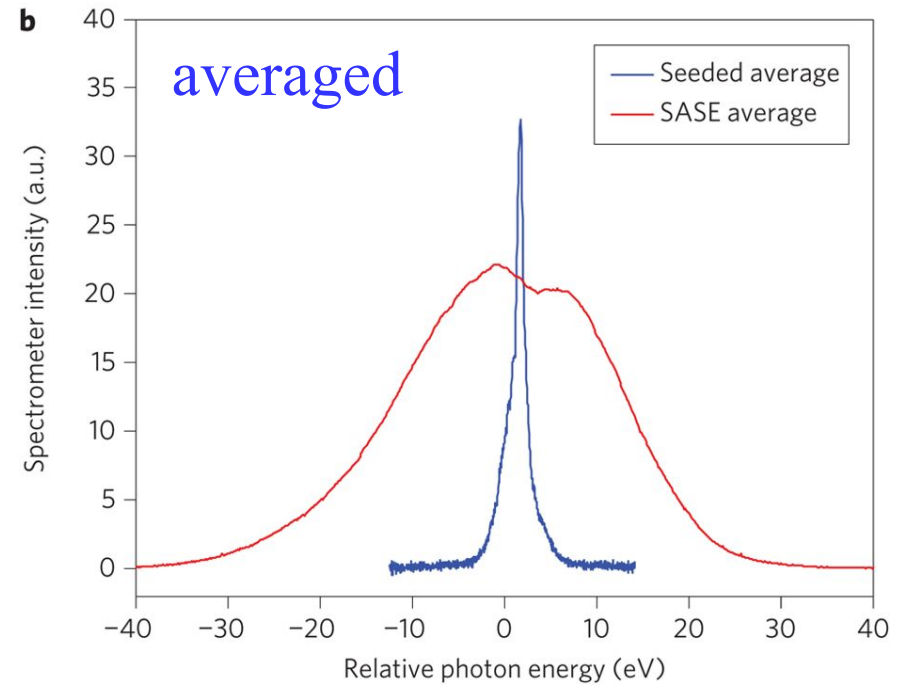
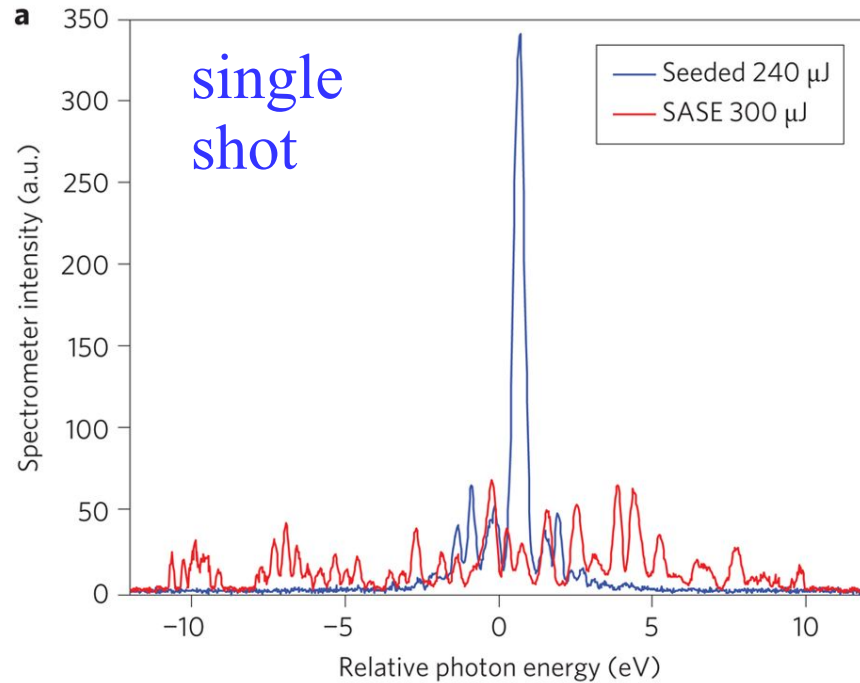


New Scientific Capabilities at European XFEL (2019)



Longitudinal coherence

Typical association:



Amann et al., Nat. Phot. 6, 693 (2012)

Relevant technical concepts:
Seeding and x-ray oscillator (XFEL)

Longitudinal coherence in more detail

“Obvious” implications:

- ▶ high spectral photon density
→ qualitatively new parameter regimes
- ▶ pulse-to-pulse stability
- ▶ reduction of heat load by off-resonant pulse components

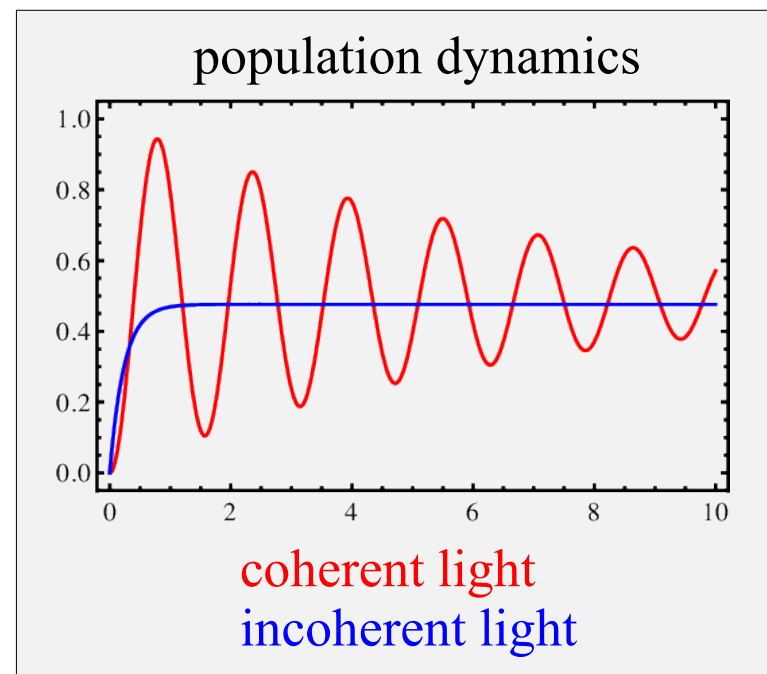
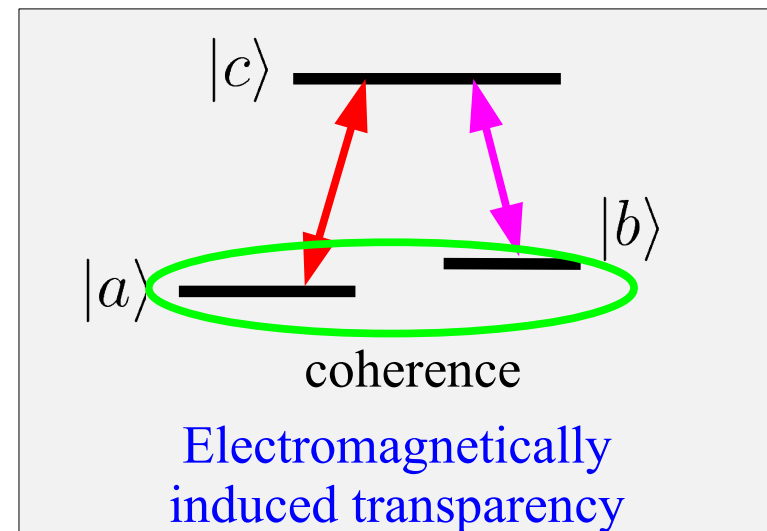
“Less obvious” implications:

- ▶ Longitudinal coherence affects the quantum dynamics of the target
→ possibility to affect/control quantum dynamics
→ advanced measurement / spectroscopy schemes
- ▶ Longitudinal coherence may enhance detection capabilities
→ interference in diffraction limited by longitudinal coherence
→ phase-sensitive “homodyne” measurements
→ Ramsey / multidimensional spectroscopy



Examples

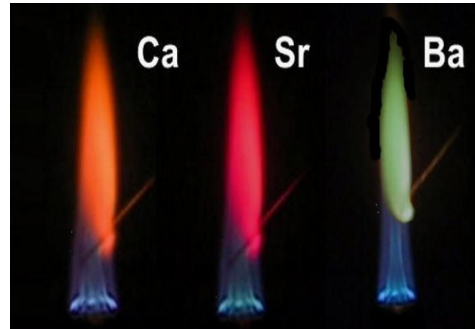
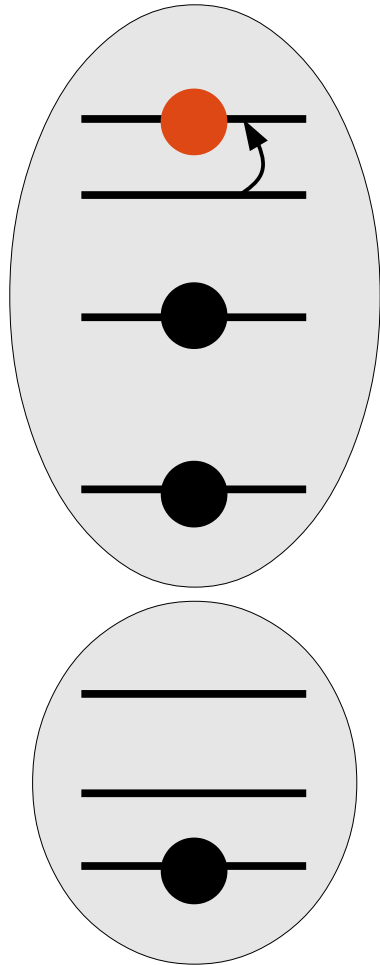
- ▶ EIT: Coherent laser fields create atomic coherence, which in turn modifies the interaction with the light
- ▶ Population dynamics: Rabi oscillation vs. incoherent rate dynamics
- ▶ Longitudinally coherent pulses yields stronger excitation than corresponding incoherent pulses



longitudinally coherent light
imprints coherence onto matter
which favorably modifies the dynamics

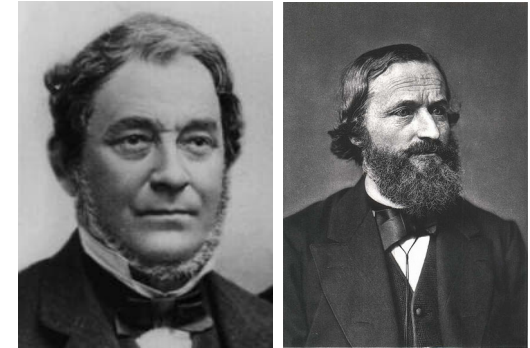
160 years of light-matter interaction in one slide

Atomic physics and quantum optics



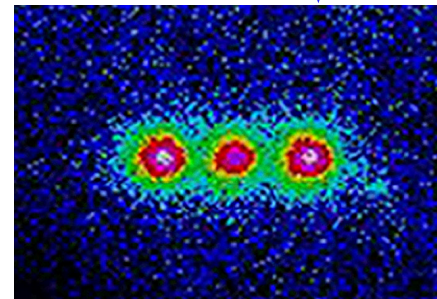
1859

“incoherent pump
and passive observation”



Bunsen Kirchhoff

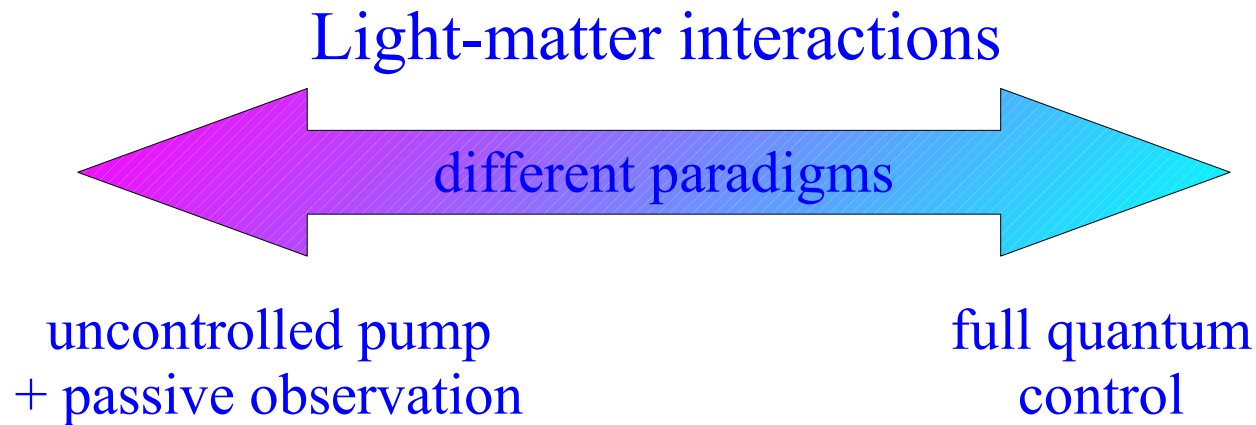
today



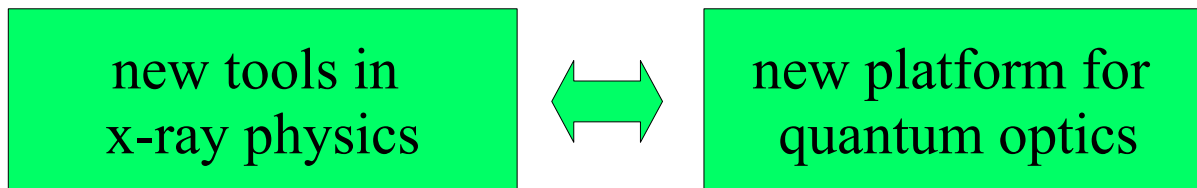
“full quantum control”

Progress enabled via coherence, non-linearities, quantum effects
→ quantum optics

X-ray quantum optics



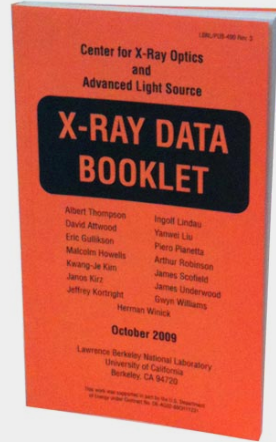
- ▶ X-ray physics could greatly benefit from moving more towards coherence/non-linear/quantum/control
- ▶ New light sources and upgrades
→ now is the right time
- ▶ Necessary to fully exploit new light sources



Not entirely new:
many existing x-ray
setups already rely
on quantum optical
concepts

Two branches of x-ray quantum optics

Electronic resonances (K-edge in ^{57}Fe)



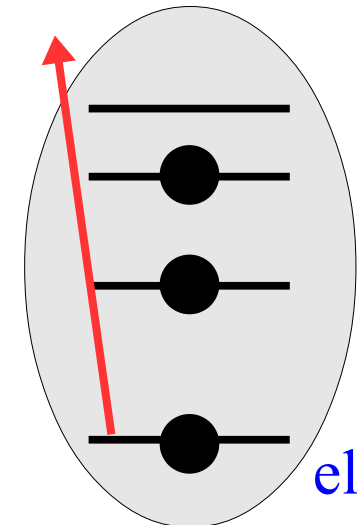
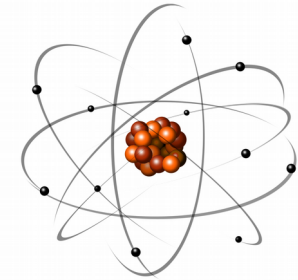
$$E = 7.112 \text{ keV}$$

$$\sigma_e = 26 \times 10^{-20} \text{ cm}^2$$

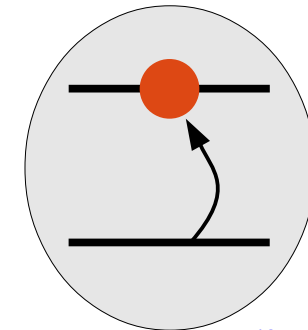
$$\Gamma \approx 2 \text{ eV}$$

$$\tau \approx 0.33 \text{ fs}$$

focus of XFEL
research

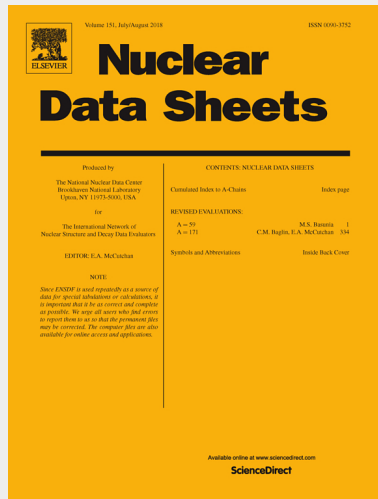


electron
shells



nucleus

Nuclear resonances (Mössbauer transition in ^{57}Fe)



$$E = 14.412 \text{ keV}$$

$$\sigma_e = 256 \times 10^{-20} \text{ cm}^2$$

$$\Gamma = 4.7 \times 10^{-9} \text{ eV}$$

$$\tau = 141 \text{ ns}$$

Extremely
narrow resonance



How could XFEL benefit from narrow resonances?

Extreme
monochromatization

Precision spectroscopy
and fundamental tests

Quantum optics
and nonlinear science

coherent
control
techniques

Narrow
resonances
as a tool

Qualitatively
new
Mössbauer
science

X-ray
quantum
optics

advanced
spectroscopy
techniques

Coherent bridge from
 \sim fs scales to
 \sim 100ns scales

Bridge between
x-ray and visible

Correlations and
(out-of-equilibrium)
dynamics

Complementary to science with electronic resonances

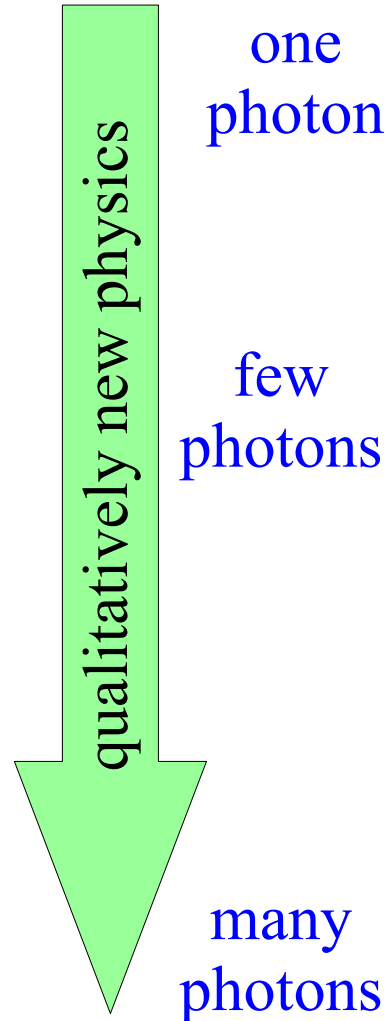


Mössbauer: Qualitatively new parameter regimes

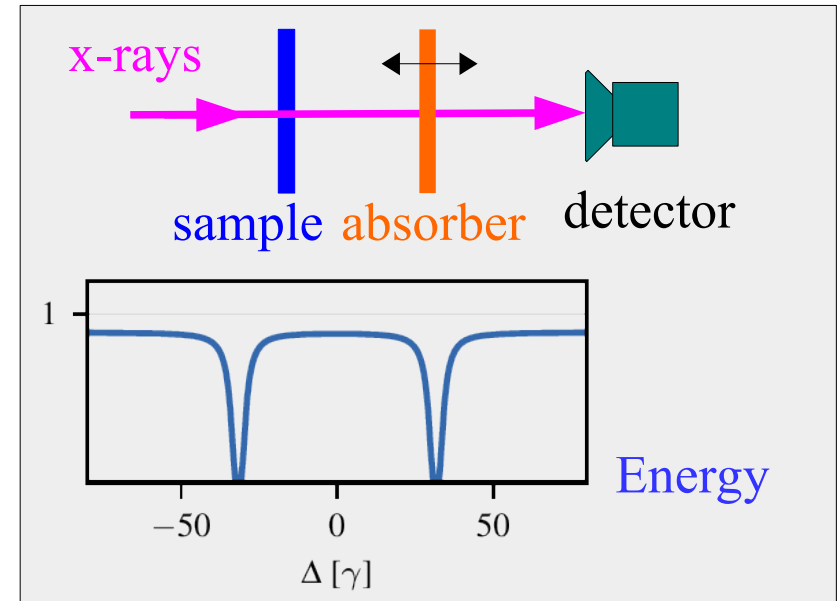
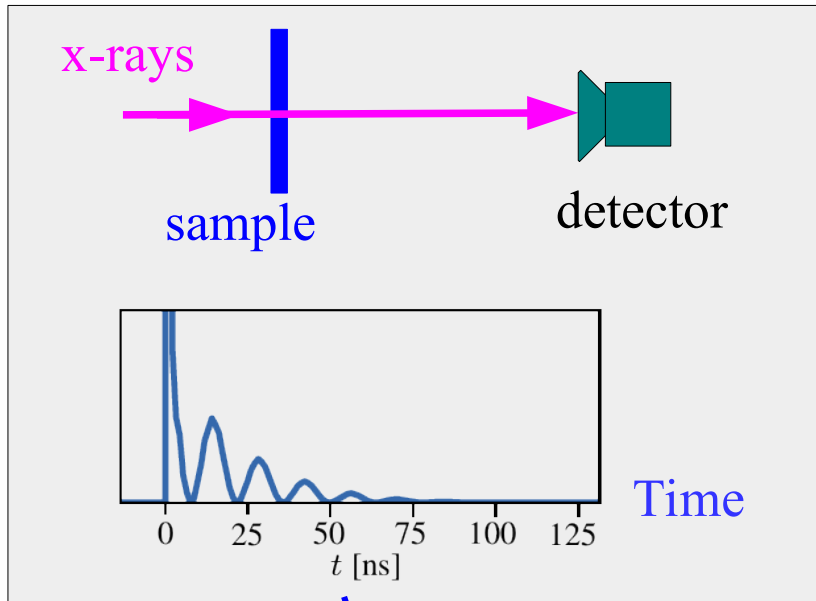
Example: ^{57}Fe

$$E = 14.4 \text{ keV} \quad \gamma = 4.7 \text{ neV}$$
$$\tau = 141 \text{ ns}$$

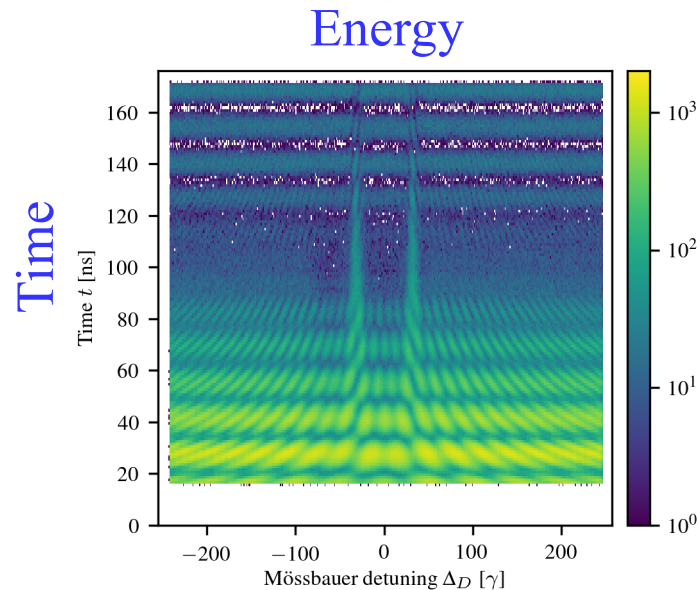
- ▶ P01 at Petra III:
on average <1 resonant photon per pulse
- ▶ SASE XFEL: (10^{11} photons/pulse, $\Delta E/E \sim 10^{-3}$)
on average ~ 30 resonant photons per pulse
- ▶ Self-seeding addition: (spectral brightness * 10)
on average ~ 300 resonant photons per pulse
- ▶ XFELo: (2.2 mJ in 28meV at 12 keV, extrapolated to 14.4 keV)
on average $\sim 10^5$ resonant photons per pulse
potentially pulse-to-pulse coherence if stabilized



Exploiting correlations – single photon case



correlate
detection
at synchrotron



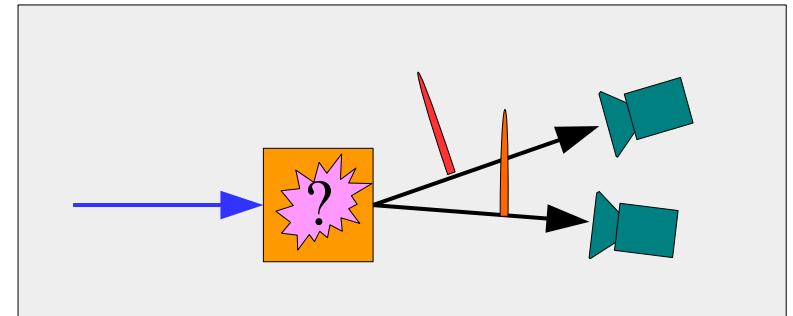
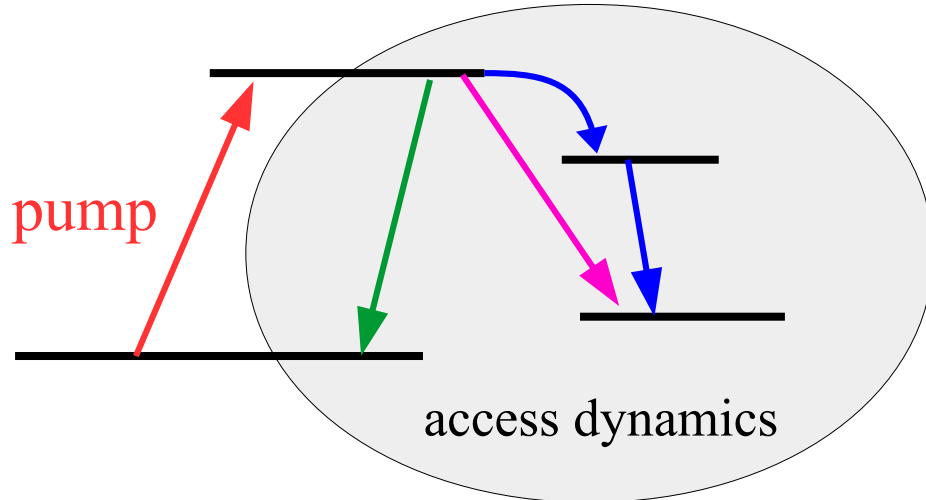
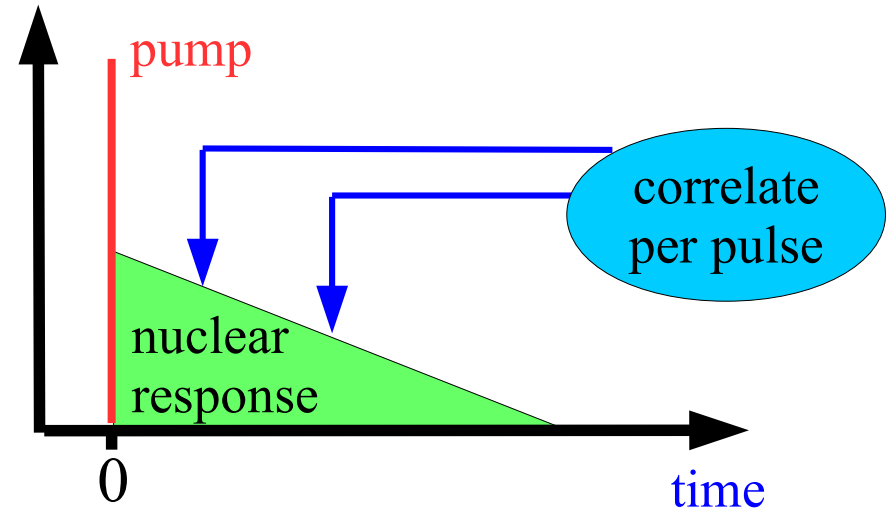
rich interference
structure provides
unique insight
into dynamics

Exploiting correlations – multi-photon case

Few signal photons per pulse

- ▶ Study pump-induced dynamics using the delayed response
- ▶ Higher-order correlations characterizing the dynamics
- ▶ Distinguish different dynamics

Intensity (log)



spatial correlations?

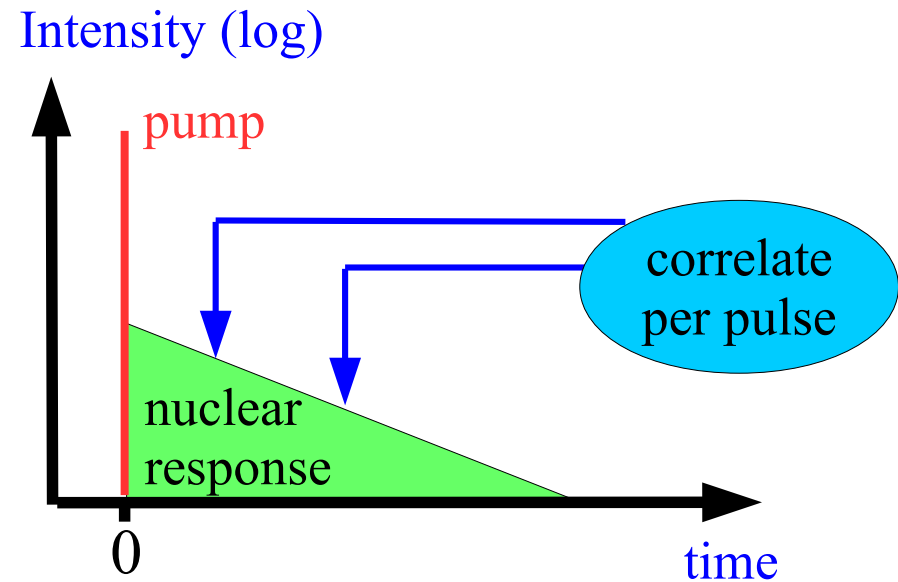
different final states

different pathways

Exploiting correlations – many photon case

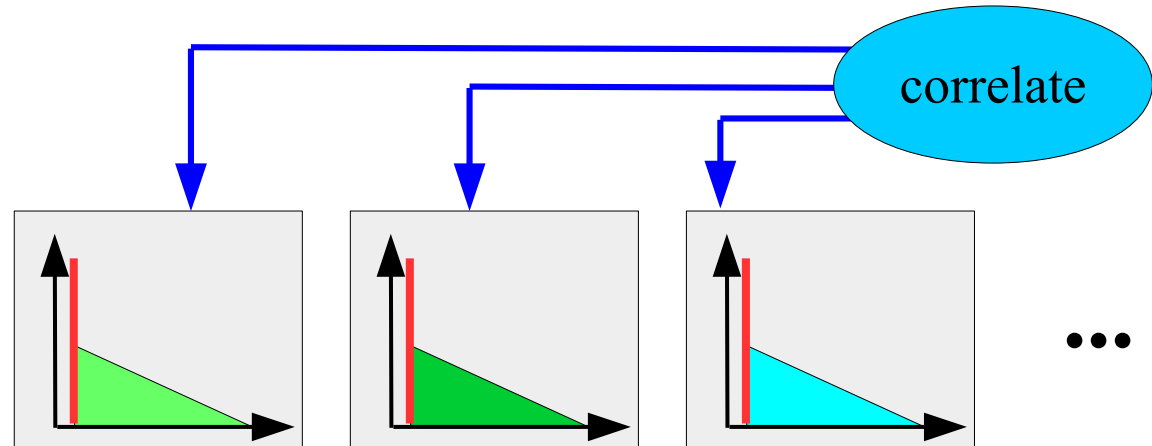
Few signal photons per pulse

- ▶ Study pump-induced dynamics using the delayed response
- ▶ Higher-order correlations characterizing the dynamics
- ▶ Distinguish different dynamics



Many signal photons per pulse

- ▶ “single shot spectra”
- ▶ compare/correlate different repetitions of pump-probe scheme
- ▶ out-of-equilibrium / non-cyclic/ non-ergodic dynamics



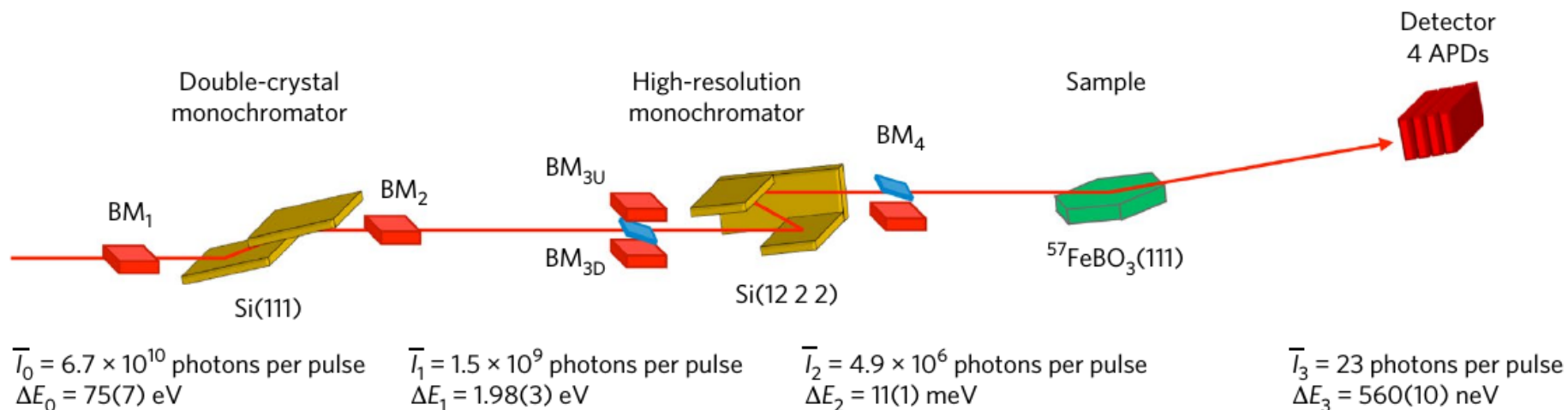
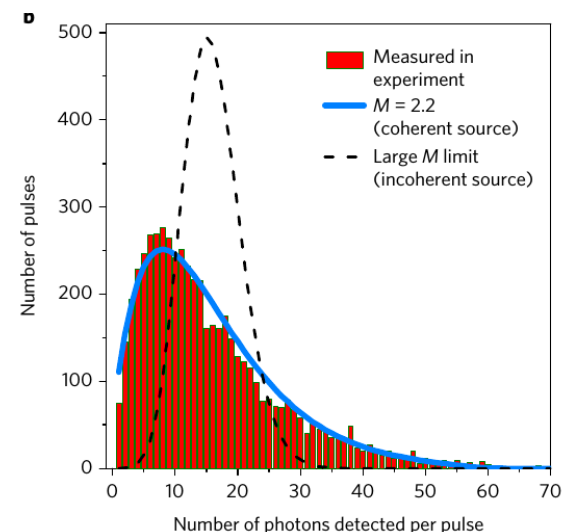
Proof-of-principle experiment

Superradiance of an ensemble of nuclei excited by a free electron laser

Aleksandr I. Chumakov , Alfred Q. R. Baron , Ilya Sergueev, Cornelius Strohm, Olaf Leupold, Yuri Shvyd'ko, Gennadi V. Smirnov, Rudolf Ruffer, Yuichi Inubushi, Makina Yabashi, Kensuke Tono, Togo Kudo & Tetsuya Ishikawa

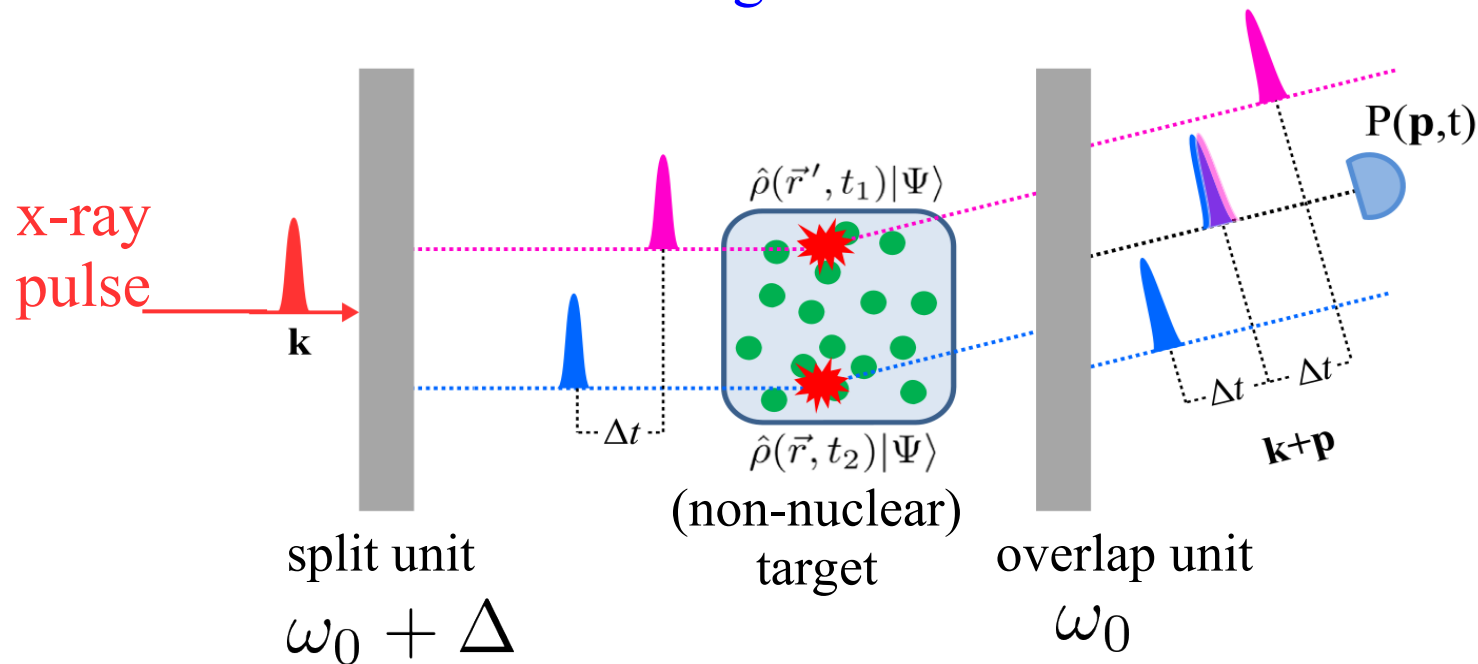
Nature Physics **14**, 261–264 (2018) | [Download Citation](#) 

- ▶ First FEL experiment with Mössbauer nuclei (at SACLA)
- ▶ Observation of correlations between photons from each shot separately, detected “one at a time”
- ▶ Example question: How does the initial emission dynamics depend on the degree of excitation?



Time domain interferometry: Nuclei as a tool

Correlations in “non-nuclear” targets:



$$P(\vec{p}, t) \propto |f(t)|^2 [S(\vec{p}, 0) + S(\vec{p}, t) \cos(\Delta t)]$$

Access to intermediate scattering function in “gap region” from ~1- ns to ~100 ns with neV energy resolution and essentially without background

(TDI proposed for applications at XFELs in SwissFEL science case)



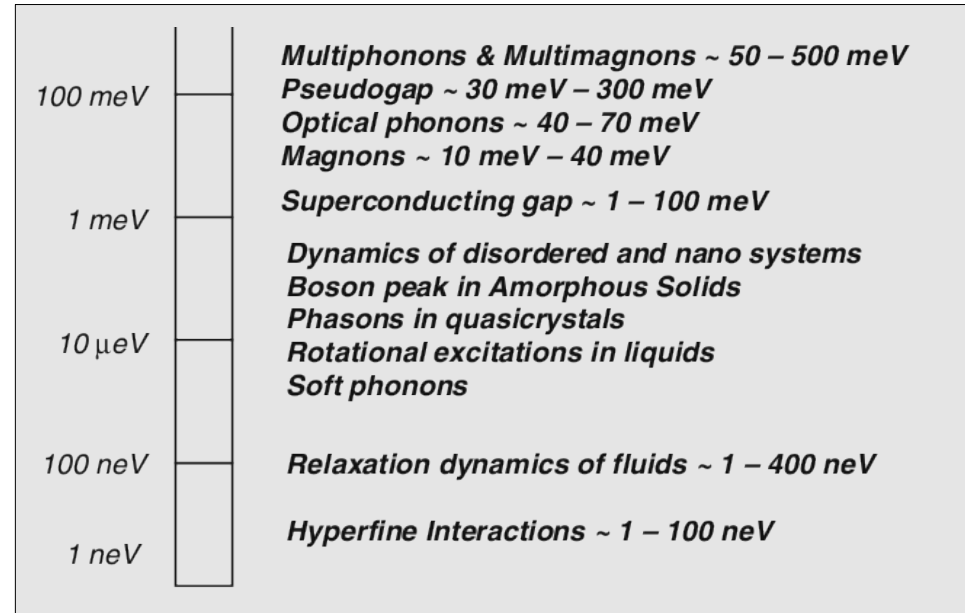
Science cases

Promote Mössbauer-based science to the study of time-dependent out-of-equilibrium phenomena

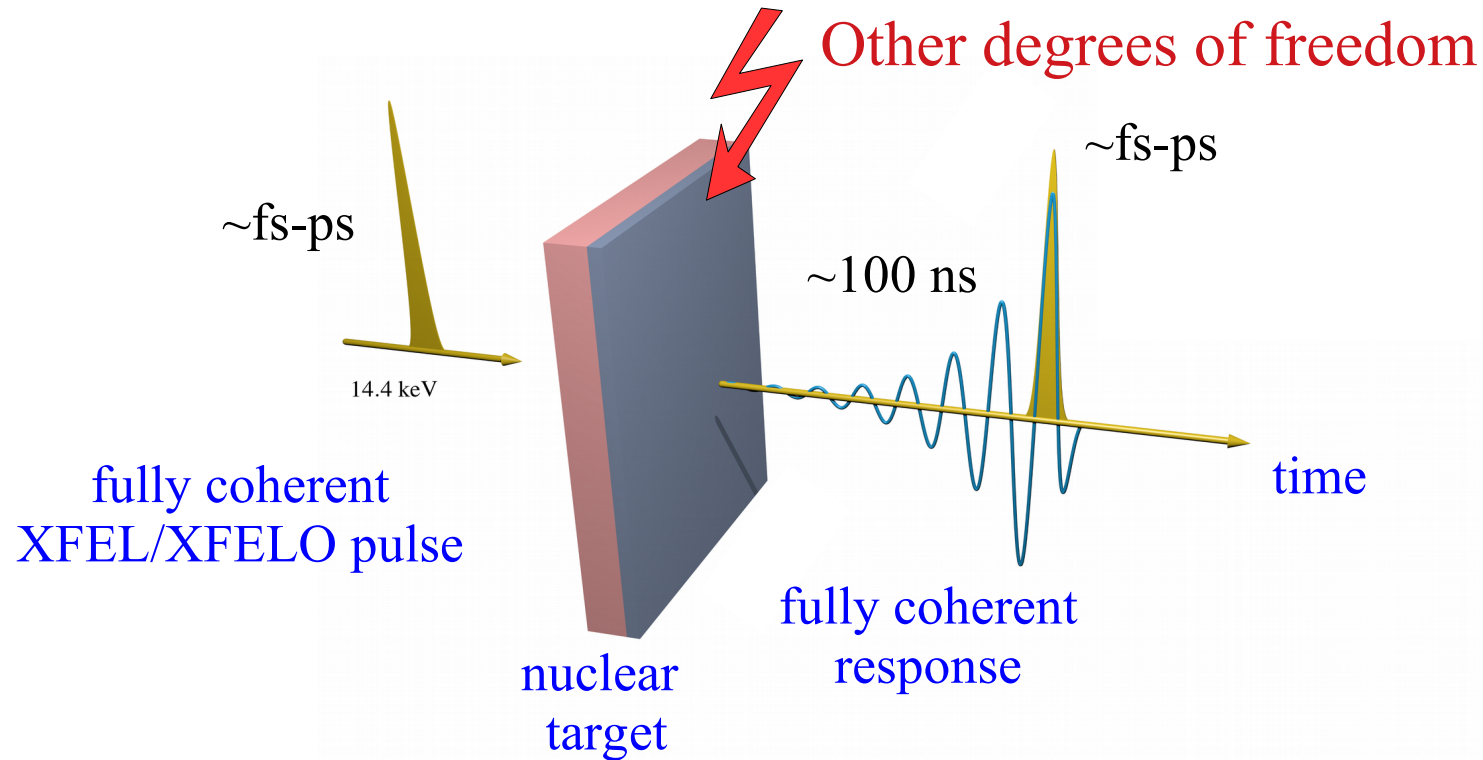
- ▶ Pump via (x-ray / optical / heat / pressure / elm. Fields / ...), probe via nuclear response
 - e.g., optically excited molecular switches
 - e.g., magneto-optical nanomaterials

Access low-energetic condensed-matter excitations on neV-meV and nm- μ m scales

- ▶ Probe low-energetic condensed-matter excitations on neV-meV and nm- μ m scales
 - e.g., physics of glasses
 - e.g., mesoscopically structured materials
 - diffusion phenomena

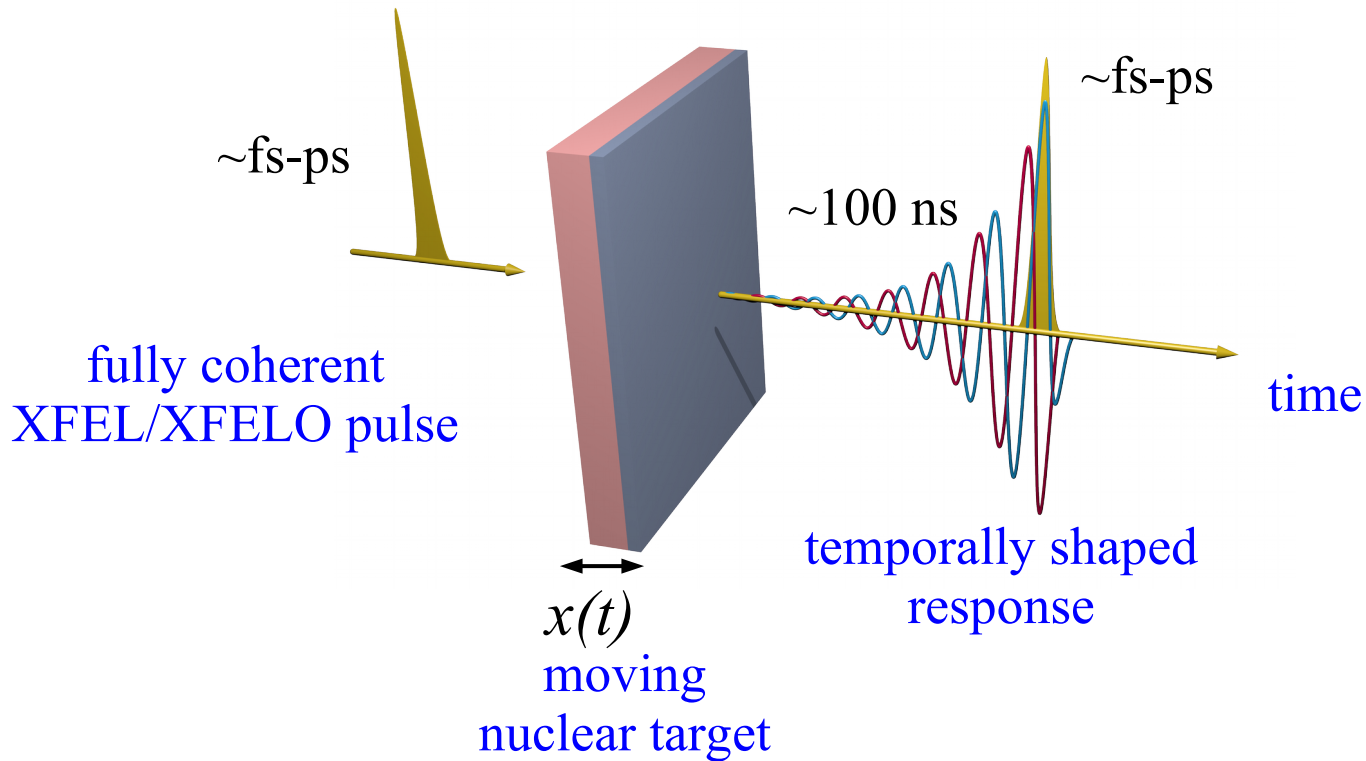


X-ray optical control of nuclei



- ▶ Nuclear resonance scattering naturally extends the longitudinal coherence to the \sim 100ns timescale
- ▶ Interference between pump pulse and response can be used for “homodyne” detection
- ▶ Can pump-probe-like ideas be used directly on nuclei?
 - probe external influence on nuclei (e.g. coupling to other subsystems)
 - important tool for nuclear quantum optics

First step: X-ray pulse shaping

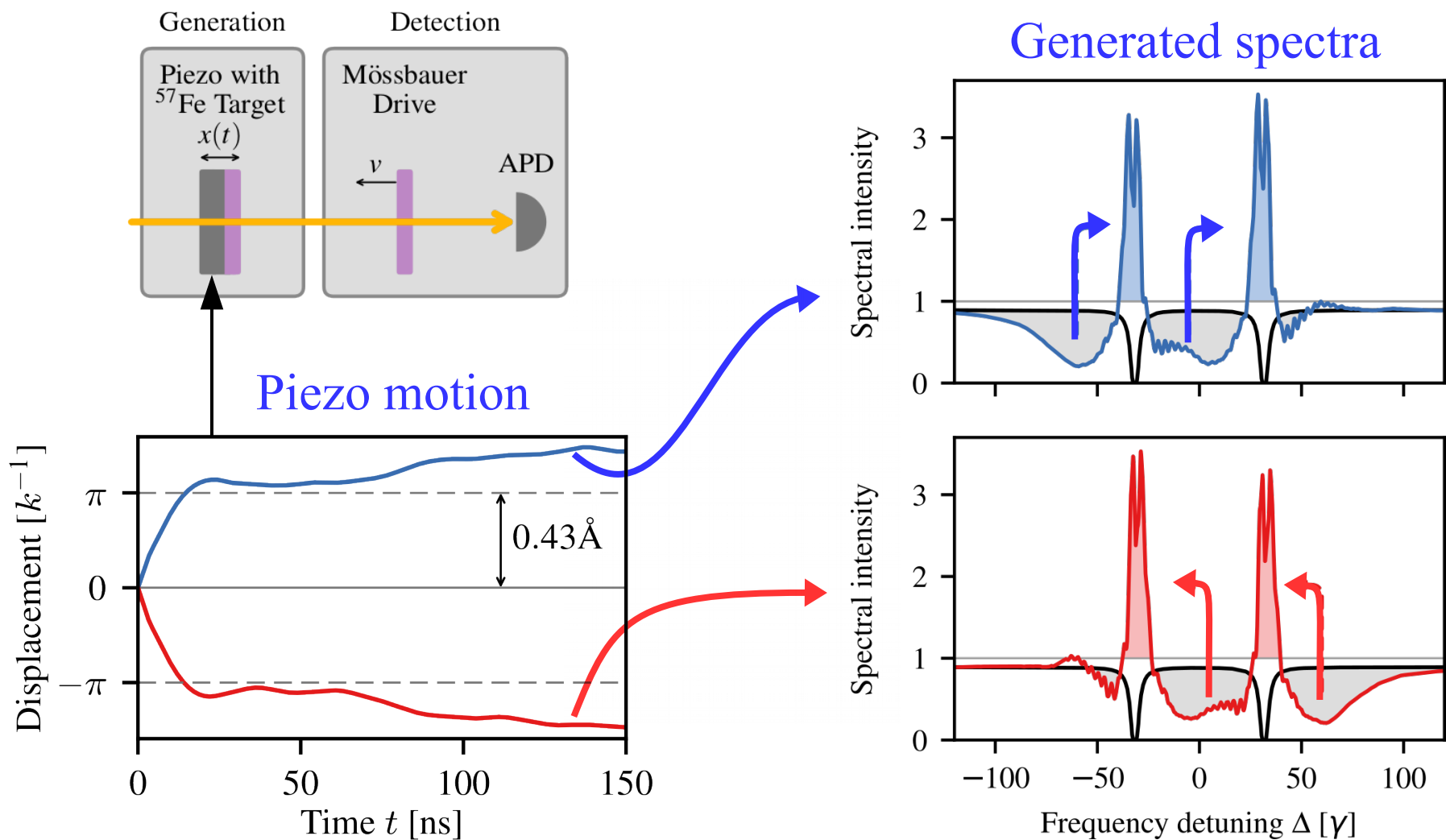


- ▶ Motion of the nuclear target imprints a time-dependent phase onto the nuclear response
- ▶ Effectively: controlled dynamical boost of real part of index of refraction by orders of magnitude without affecting absorption/imaginary part

$$n = 1 - \delta[x(t)] + i\beta$$



“X-ray afterburner” (ID 18, ESRF)

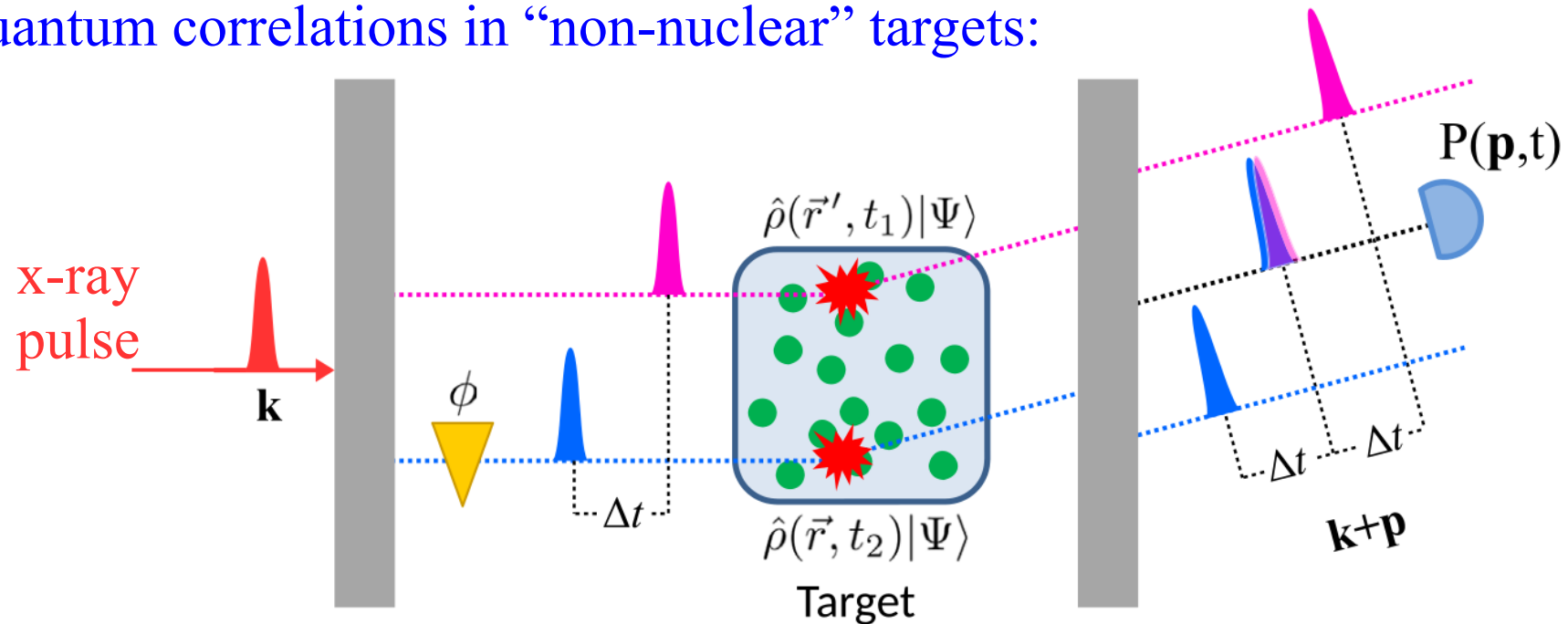


Resonant “gain” exceeding all loss channels
Mechanical motion controlled and measured on ns / sub-Å level

Black line:
no motion

Immediate application: Phase-control in TDI

Quantum correlations in “non-nuclear” targets:

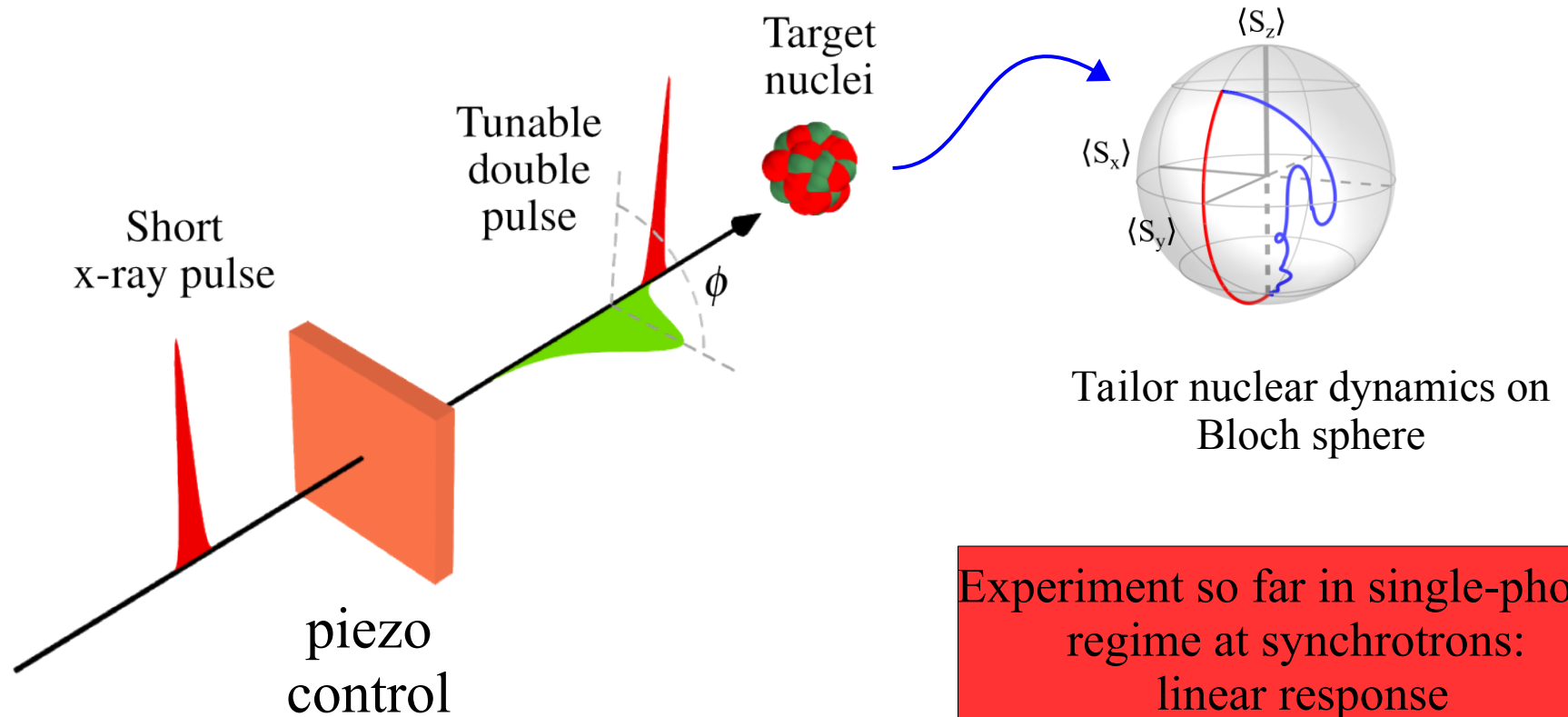


$$P(\mathbf{p}, t) \propto f(t)^2 \left(\sum_{j=1,2} S_{qu}(\mathbf{p}, t_j, t_j) + 2 \left\{ \cos[\phi] S_{qu}^R(\mathbf{p}, t_1, t_2) - \sin[\phi] S_{qu}^I(\mathbf{p}, t_1, t_2) \right\} \right)$$

real part
imaginary part

Phase-control of one of the two scattering pathways provides access to real and imaginary parts of the complex intermediate scattering function
 → quantum mechanical correlations

Towards nuclear pump-probe experiments



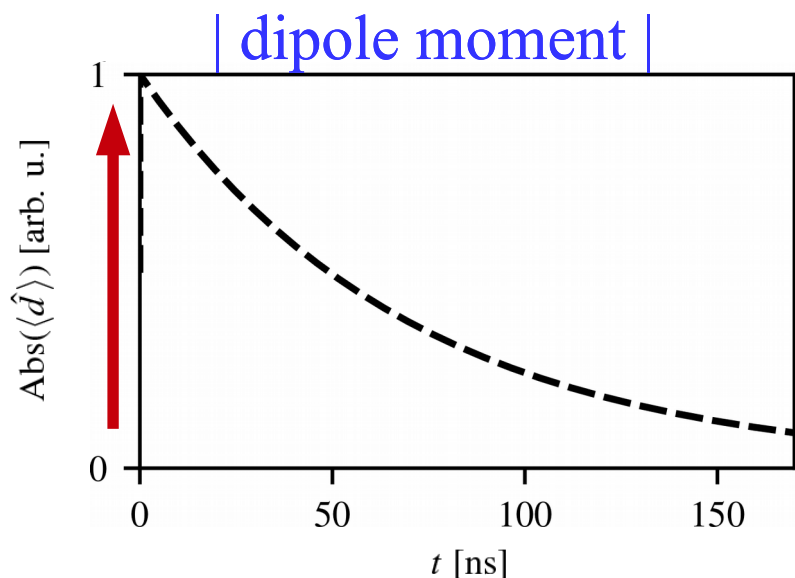
Experiment so far in single-photon regime at synchrotrons:
linear response

”true” pump-probe requires strong driving \rightarrow XFEL(O)

- ▶ First pulse excites nuclear target
- ▶ Piezo control shapes second pulse part
- ▶ Double-pulse determines the dynamics of the nuclear target

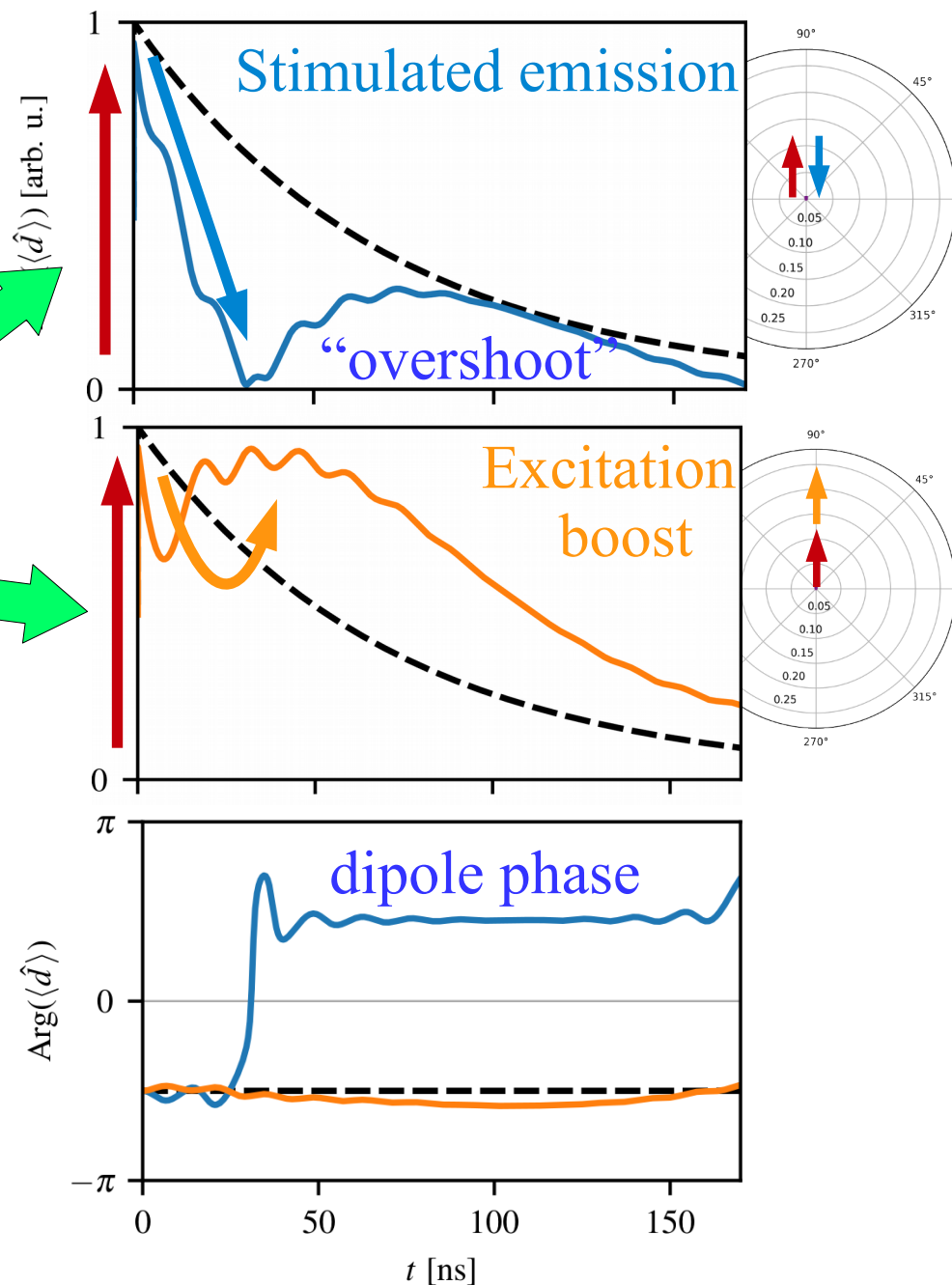
Experimental results: x-ray optical control of nuclei

Preparation pulse
excites the system:



Regular decay without
second pulse

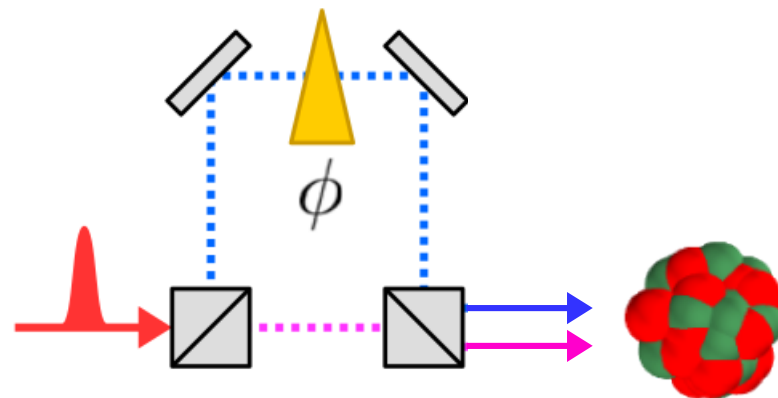
X-ray optical control of
nuclear dynamics
stable to fractions of x-ray wavelength



Why is the control so stable?

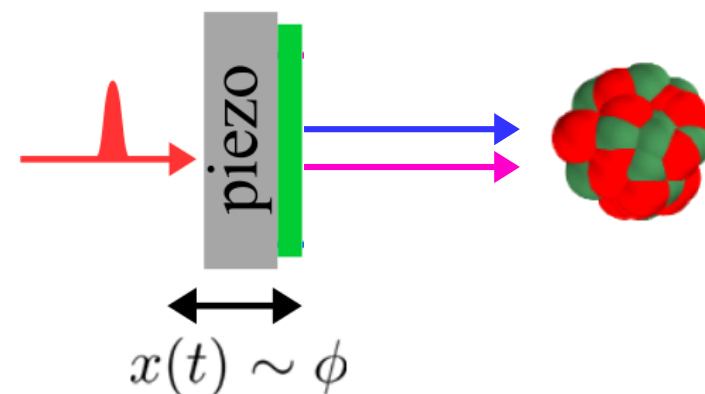
“Conventional approaches”

- ▶ Interfering pathways spatially separated
- ▶ Geometry must be stabilized throughout the entire long accumulation of statistics



Piezo control with mechanical motion

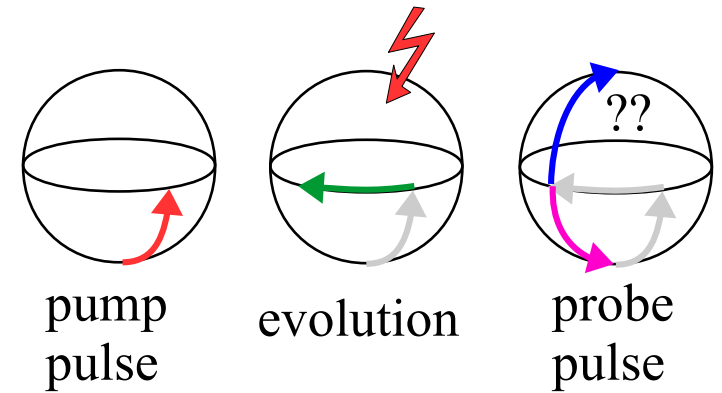
- ▶ Interfering pathways coincide in space
- ▶ Control depends on motion relative to the geometry at the time of excitation
- ▶ Geometry only needs to be stable for a ~ 200 ns measurement interval after each x-ray pulse
- ▶ All other drifts / noise do not matter!



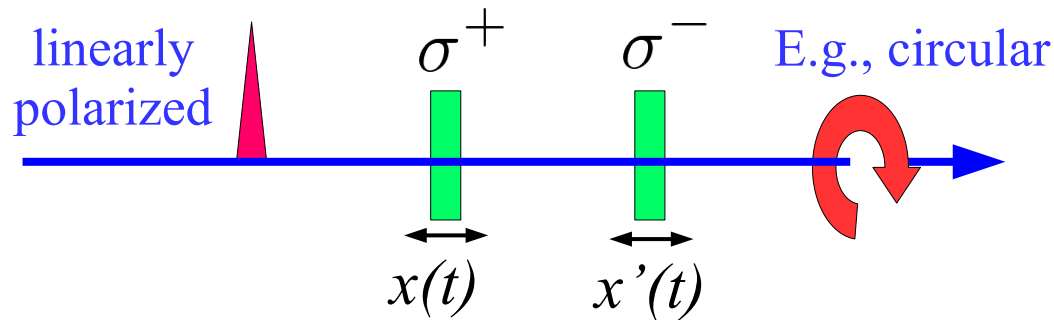
Piezo control for nuclear quantum optics

Advanced spectroscopy (e.g., Ramsey)

- ▶ Robust method to precisely measure frequencies, dynamics of coherences, external couplings
- ▶ First step towards multidimensional spectroscopy

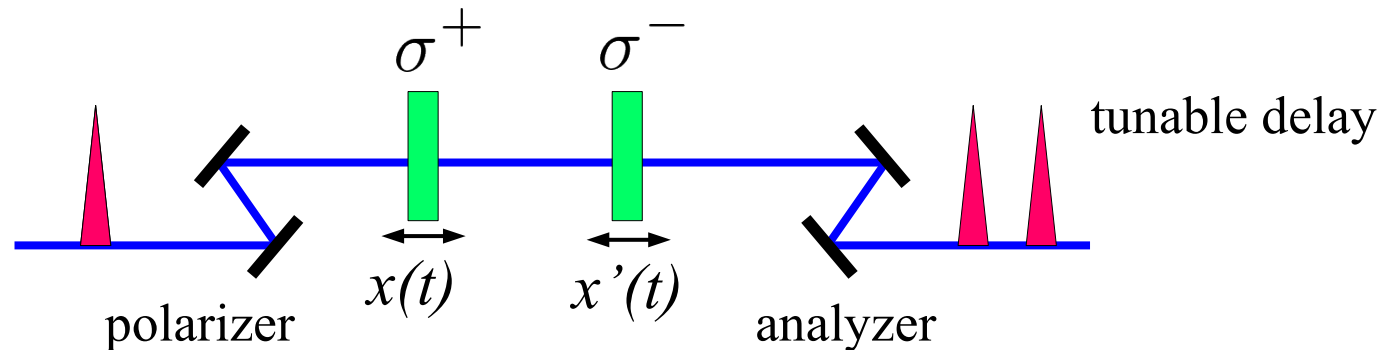


Dynamical polarization control



Motion-induced birefringence/ "waveplates"

"True" double pulses



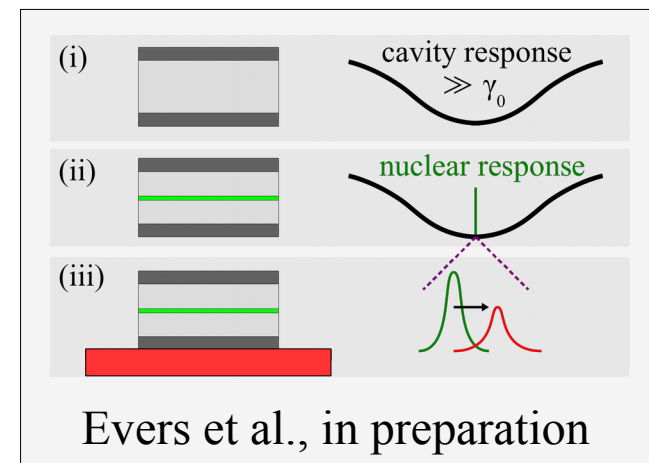
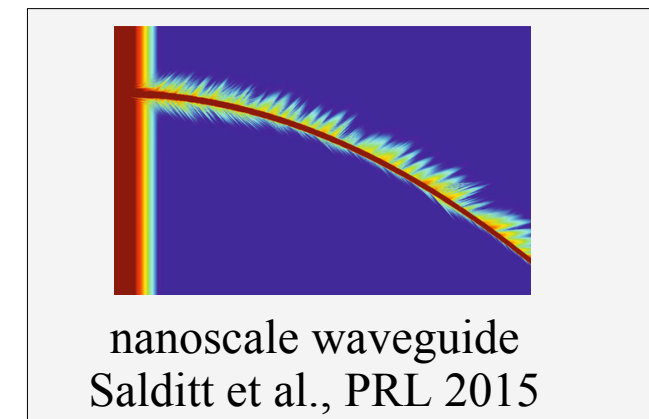
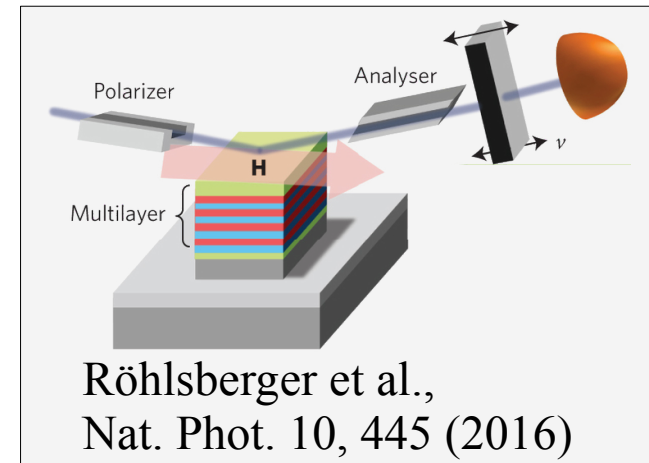
Crucial: Target design

▶ Smaller targets / higher resolution

- interfaces (magnetism, friction, ...)
- spatially resolved dynamics (heat, magnetization, chemistry)

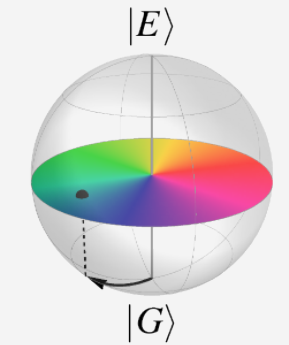
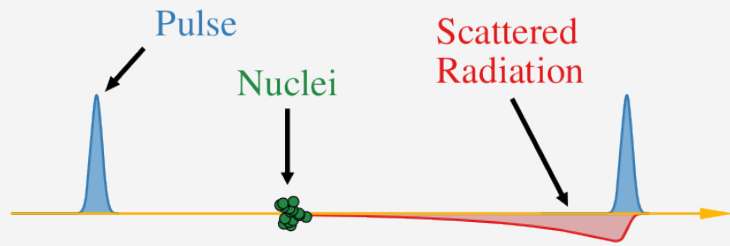
▶ Structured targets

- design artificial quantum systems (e.g., EIT with two-level systems)
- enhance interaction between light and matter (e.g., via cavities, waveguides)
- x-ray nanophotonics: photonic crystals / light-matter quasiparticles / x-ray photonic circuits
- decoherence control via engineering of environment (e.g. SGC)
- evanescently coupled sensing schemes

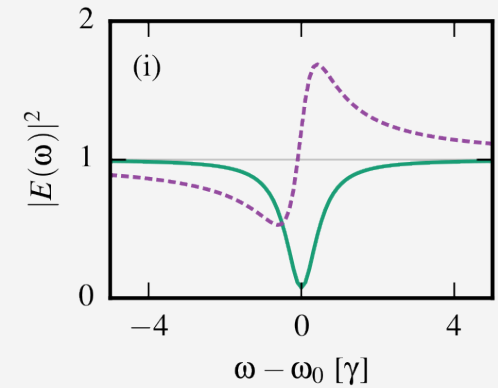


Towards non-linear science with nuclei: Rabi flopping

Weak x-ray pulse

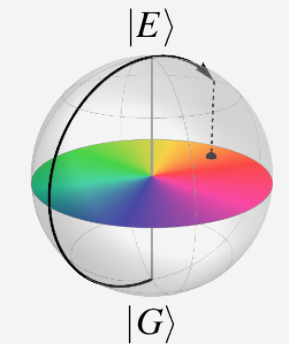
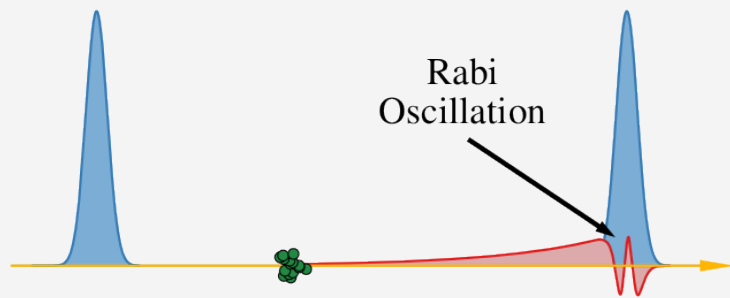


Bloch sphere

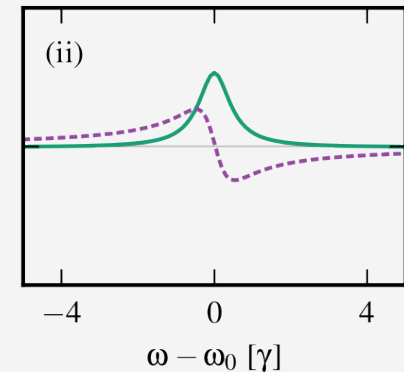


Observed spectrum

Strong x-ray pulse



Bloch sphere



Observed spectrum

Dipole phase sign change upon each half Rabi cycle leads to flip of spectra

$$E(t) = A_0 \delta(t) - \frac{\beta d}{2} \sin(\Phi) \theta(t) e^{-\frac{\gamma}{2} t}$$

Nonlinear nuclear quantum optics

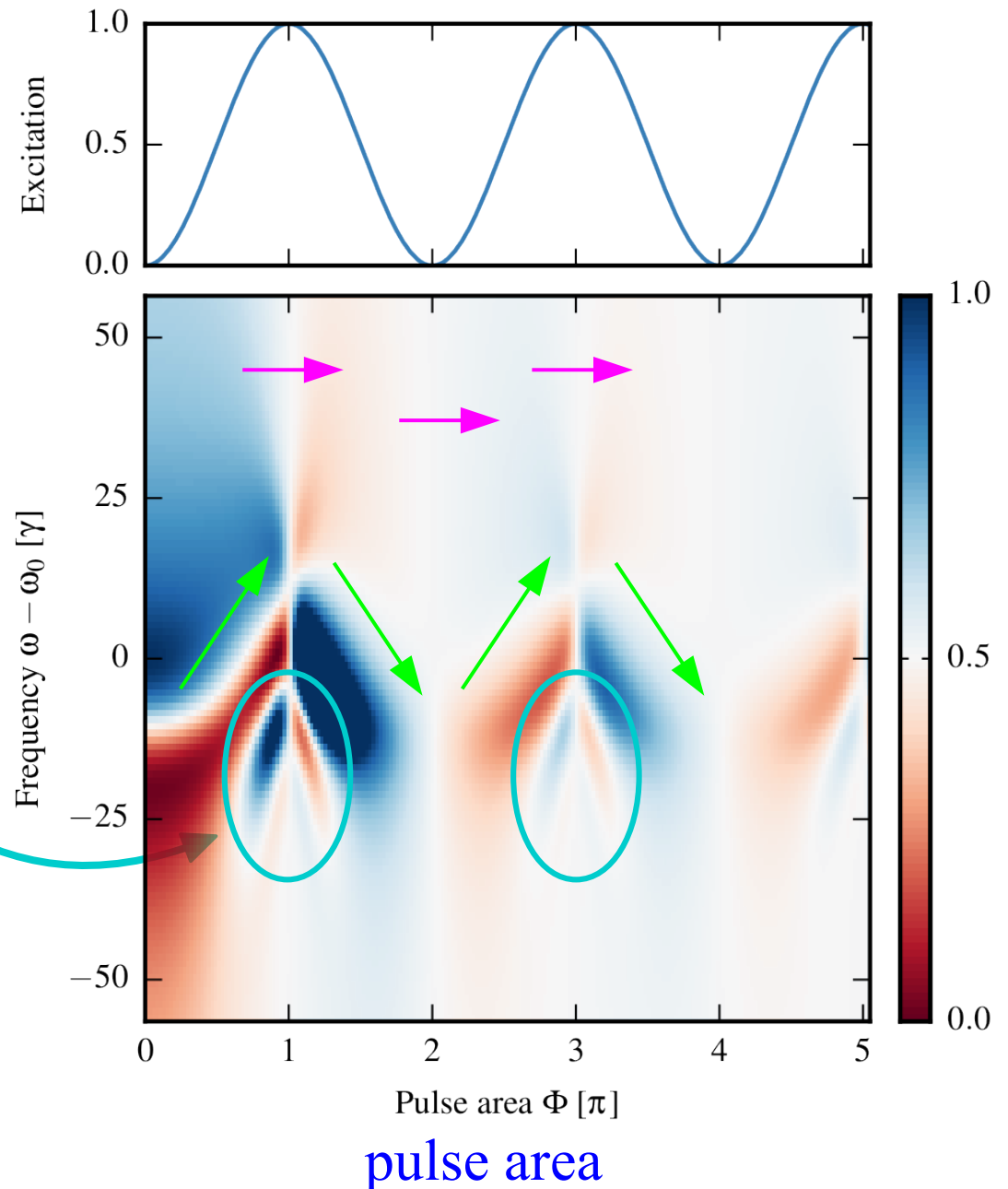
Signatures of strong excitation of nuclei in x-ray cavities

- ▶ Symmetry flips of the spectra of the delayed photons
- ▶ Corrections due to collective Lamb shifts
- ▶ Additional spectral signatures due to non-exponential collective decay

Non-linear signatures may be within reach with XFELs, focusing, optimized sample design

also c.f. Chumakov FEL experiment

spectrum



Quantum interfaces linking x-rays and “low frequency”

Vision: **coherent** link between x-rays and lower frequencies to combine best of both “worlds”

$4.7 \text{ neV} / \hbar = 1.1 \text{ MHz}$ – comparable scales

► X-ray optomechanics

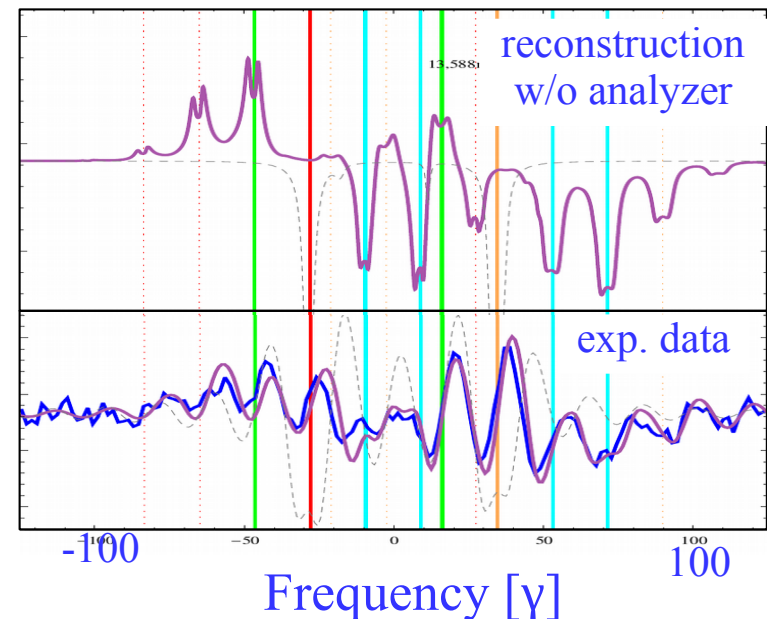
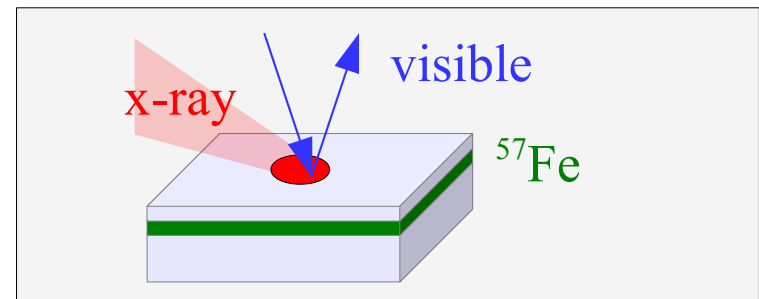
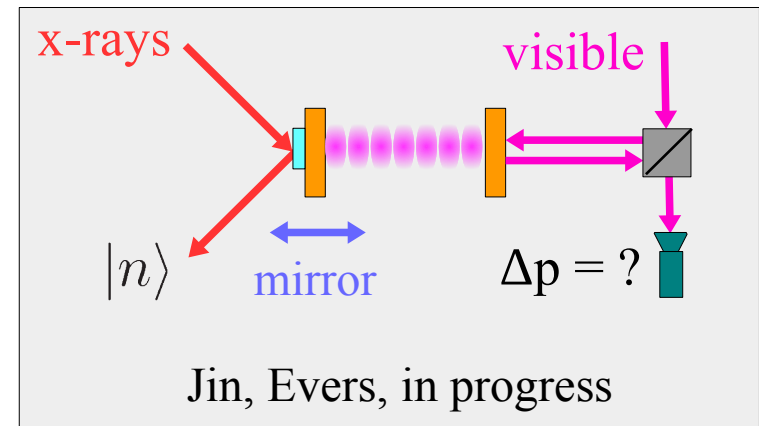
- Manipulate x-rays using optical field (proposal Palffy et al, 2016)
- Macroscopic x-ray Fock state generation

► Coupling via common target

- x-ray detection of laser-induced motion (e.g. Kocharovskaya et al)
- visible-pump / nuclear probe (e.g. Schünemann et al)

► Nuclear frequency combs

- Modulate nuclear spectra using RF-fields (higher frequencies?) (Kocharovskaya et al, Evers, Pfeifer, Röhlberger et al)



Summary

- ▶ Narrow nuclear resonances at XFELs:
 - useful as a tool
 - qualitatively new Mössbauer science
 - platform for x-ray quantum optics
- ▶ Seeding/XFEL schemes provide access to qualitatively different regimes for nuclear resonances
- ▶ Promote nuclear resonance scattering to time resolved (pump-probe) study of dynamical (out-of-equilibrium) processes
- ▶ Access to low-energetic solid-state excitations
- ▶ Coherent x-ray optical control of nuclei demonstrated in single-photon regime
 - nuclear pump-probe with seeding/XFEL?
 - Ramsey / multidimensional spectroscopy?
 - probe external influence on nuclei, other subsystems
- ▶ Interfaces between x-ray and “low frequency”

Acknowledgements:

Heeg, Kaldun, Strohm, Reiser, Ott, Subramanian, Lentrod, Haber, Wille, Goertler, Ruffer, Keitel, Röhlberger, Pfeifer, Evers, Science 357, 375 (2017)

Castrignano & Evers, PRL 122, 025301 (2019)

MPIK Heidelberg, Germany

DESY, Hamburg, Germany

ESRF, Grenoble, France

Report on the 2016 SLAC workshop on the scientific opportunities of XFELs arXiv:1903.09317 [physics.ins-det]

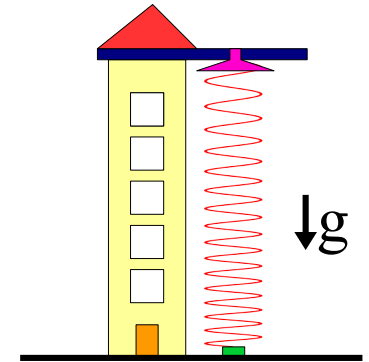
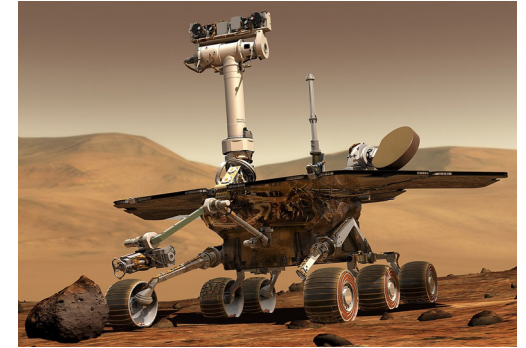
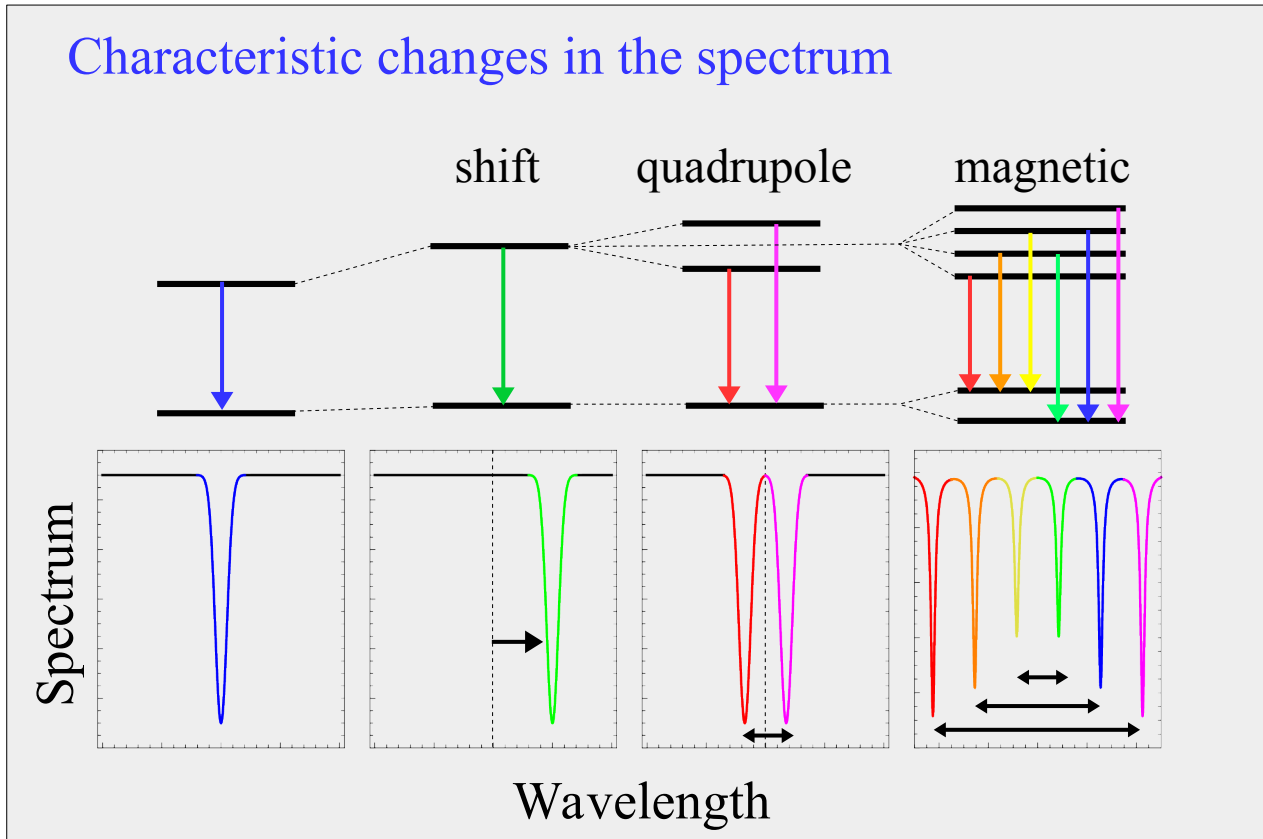
Shenoy&Röhlberger, Hyperf. Int. 182, 157 (2008)

Discussions with many colleagues





Mössbauer nuclei: sensitive probe of the environment

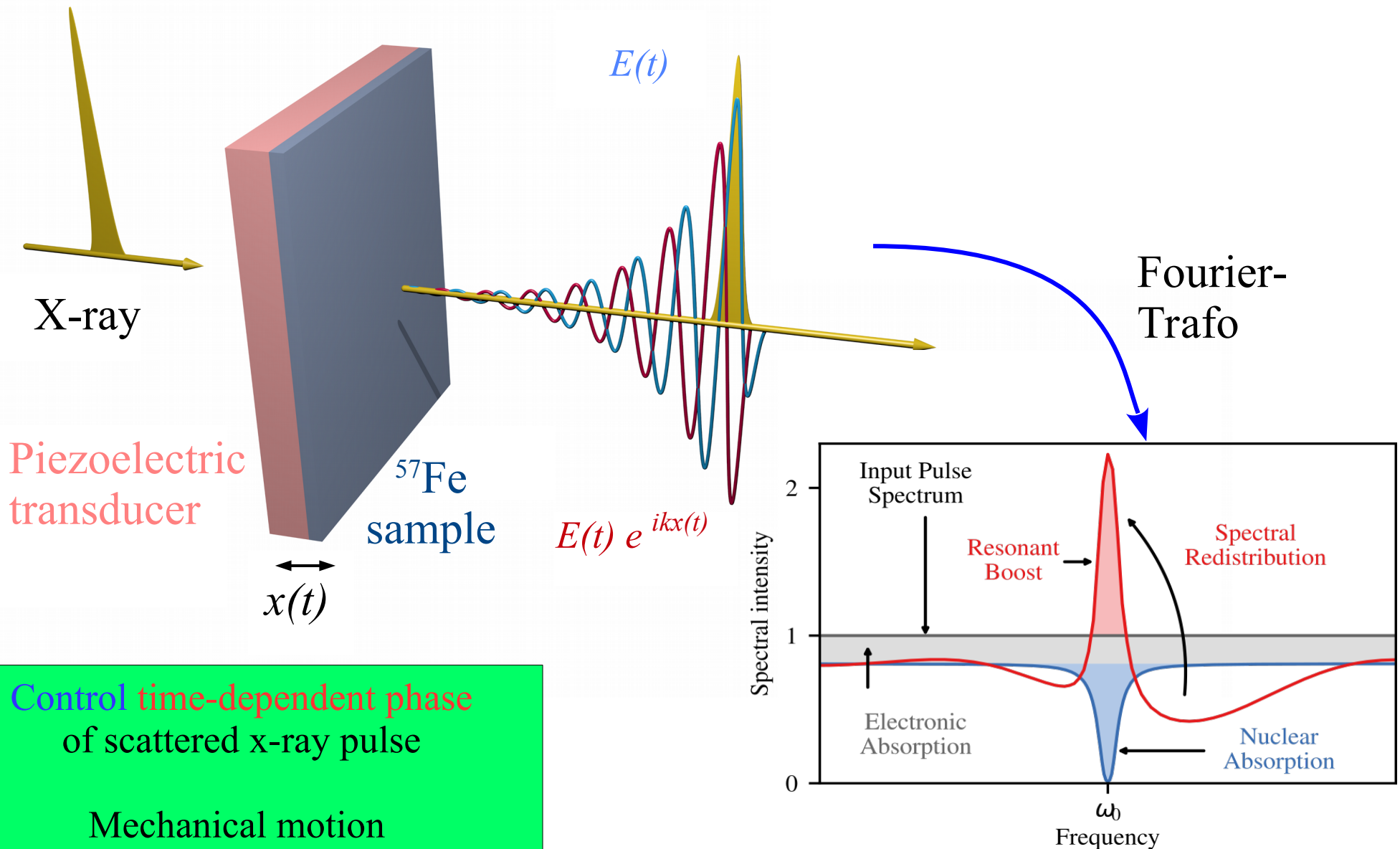


- ▶ identification of chemical, spin, and magnetic state
- ▶ relaxation phenomena
- ▶ coordination to neighbors
- ▶ value and direction of magnetic field



+ inelastic scattering: phonon density of states, dispersion relations, relaxation, ...

X-ray pulse shaping



Control time-dependent phase
of scattered x-ray pulse

Mechanical motion
simulates control laser field

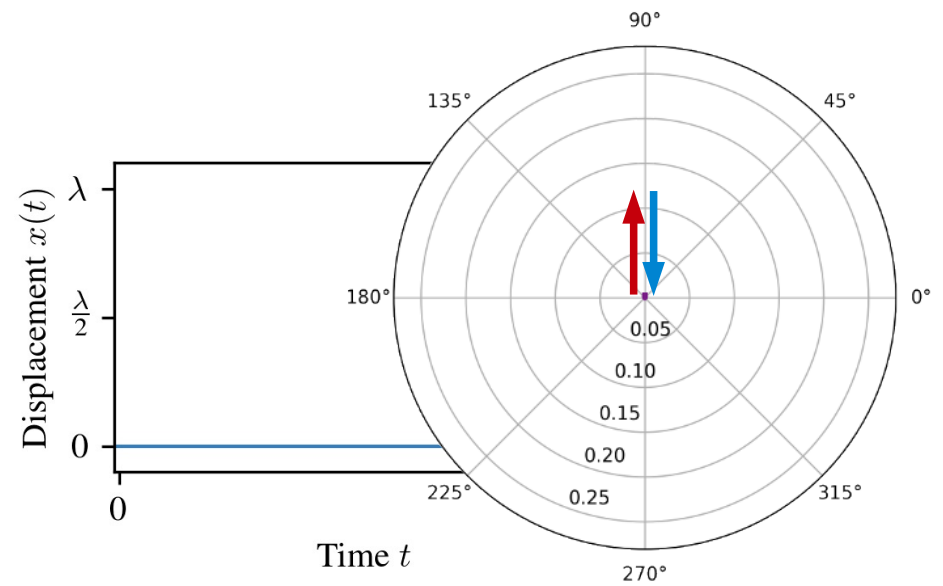
Demonstrate control with two special cases

► “Stimulated emission of excitons”

Two pulses with **opposite phase**:

First preparation pulse
excites exciton

Second control pulse
de-excites target

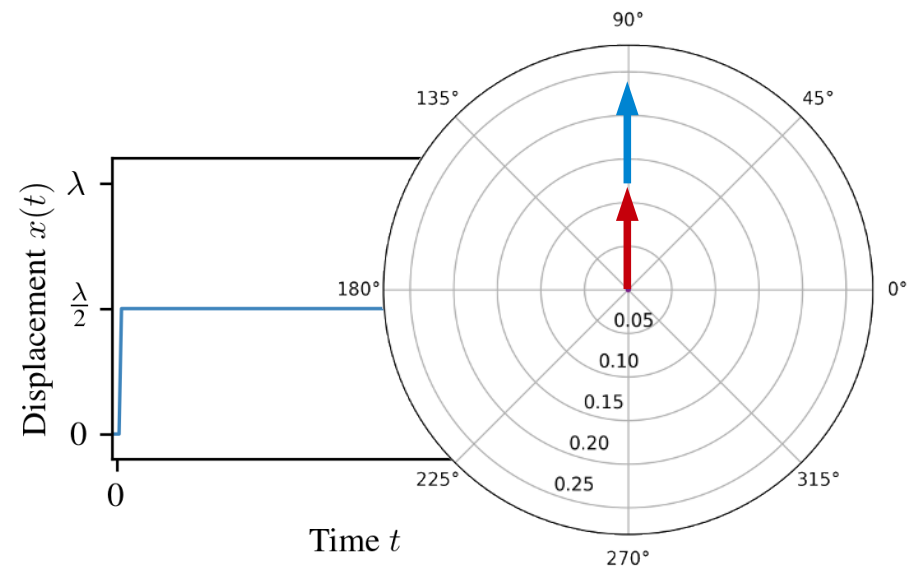


► “Coherent boost of excitation”

Two pulses with **same phase**:

First preparation pulse
excites exciton

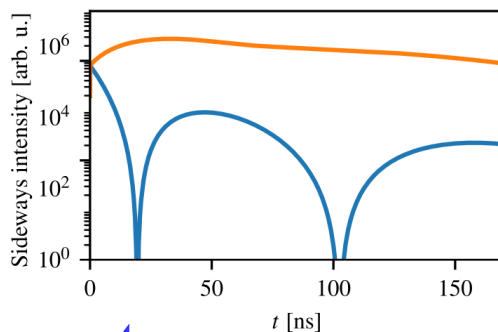
Second control pulse
further excites target



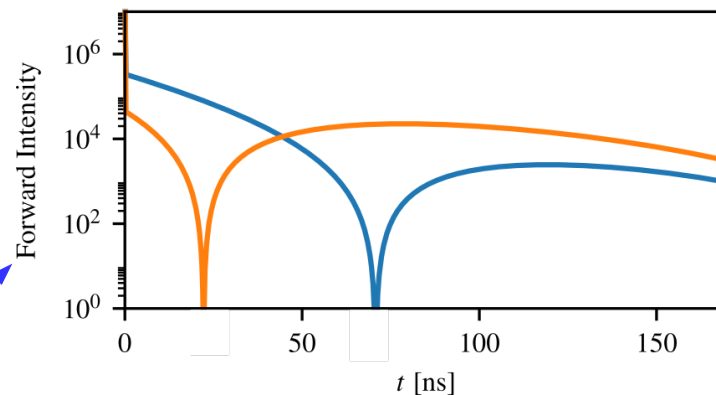
How to observe?

Sideways:

- Only scattered light
- Clear signature
- Low count rate

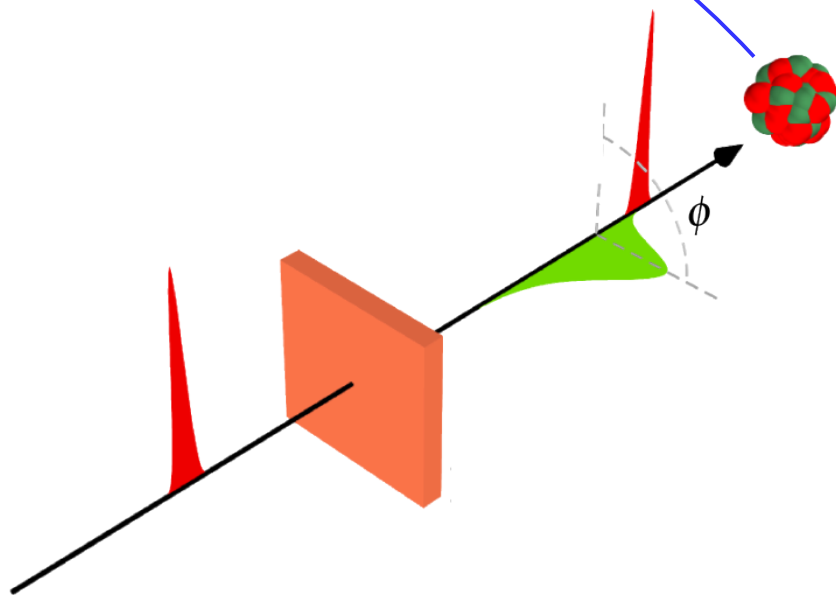


Stimulated emission Excitation boost



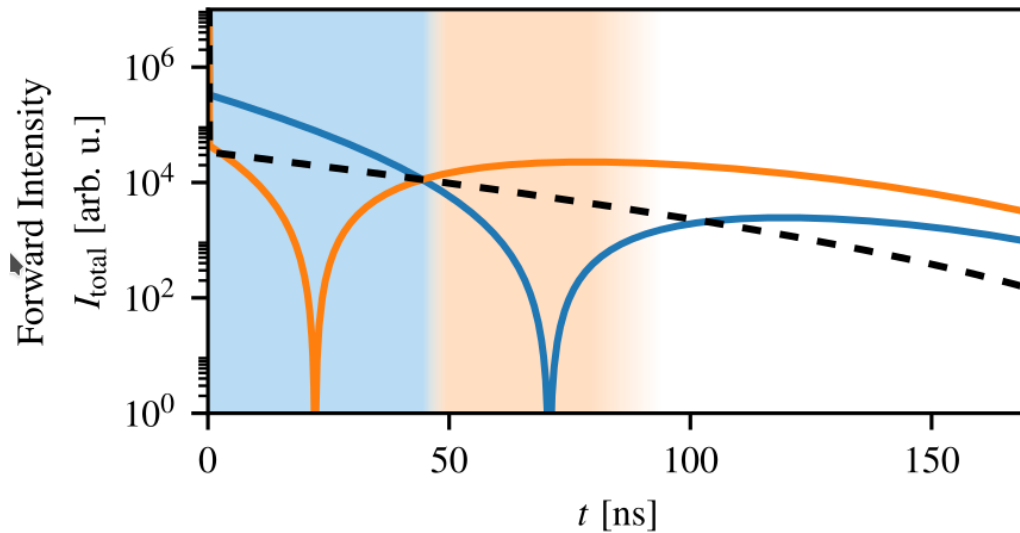
Forward direction:

- Interference with incident pulses
- Phase-sensitive
- High count rate
- Signature of dynamics?



Signature for spontaneous emission and boost

Theory model for single resonance in target



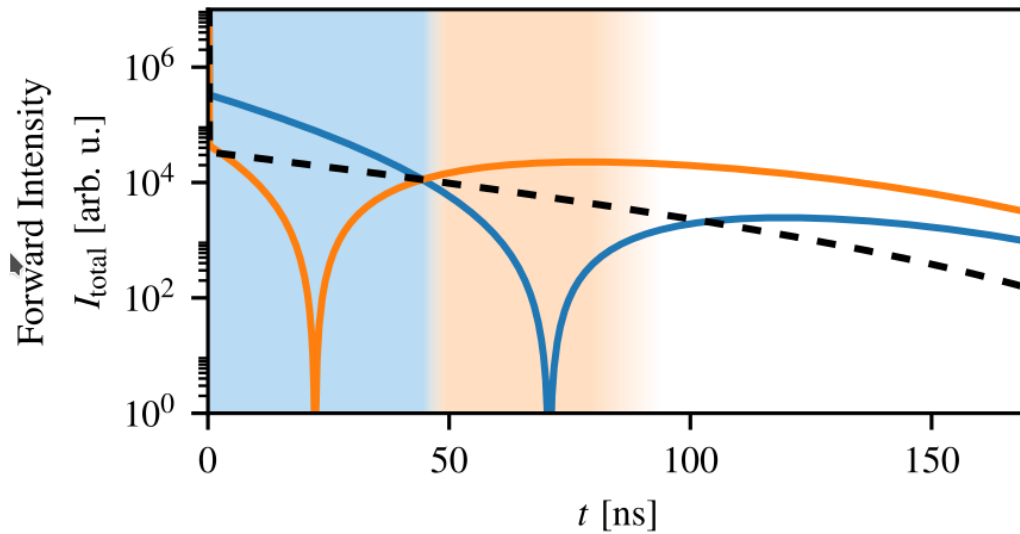
- ▶ Initially, higher intensity with SE, due to faster decay
- ▶ Afterwards, lower intensity with SE, since nuclei have already decayed

Stimulated emission
Excitation boost



Experimental results (ID 18, ESRF)

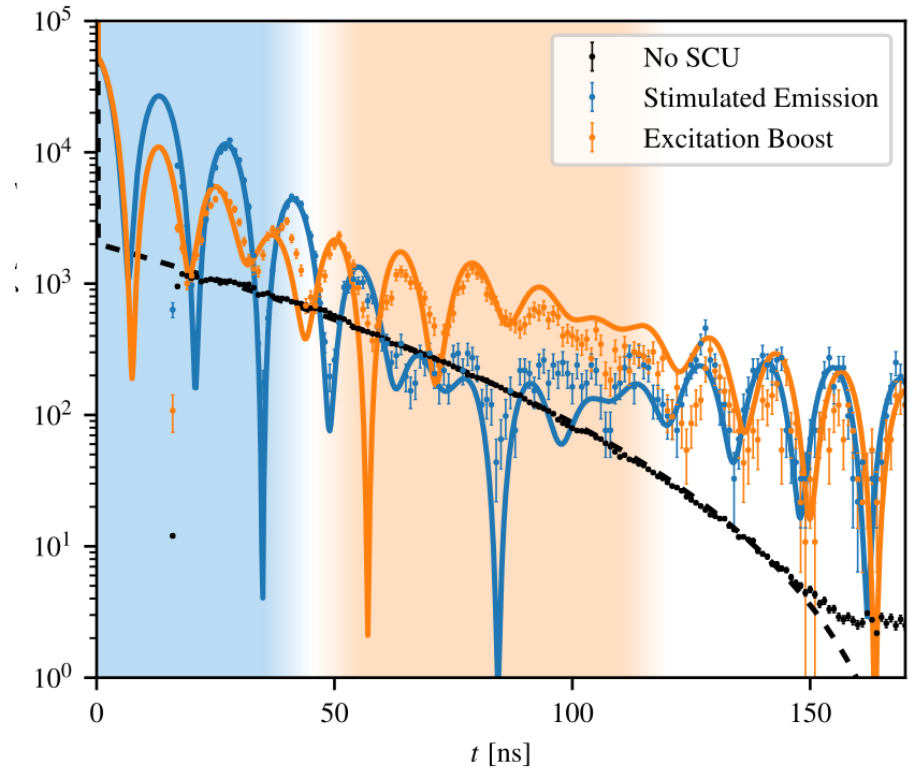
Theory model for single resonance in target



- Initially, higher intensity with SE, due to faster decay
- Afterwards, lower intensity with SE, since nuclei have already decayed

Stimulated emission
Excitation boost

Experimental results with fit



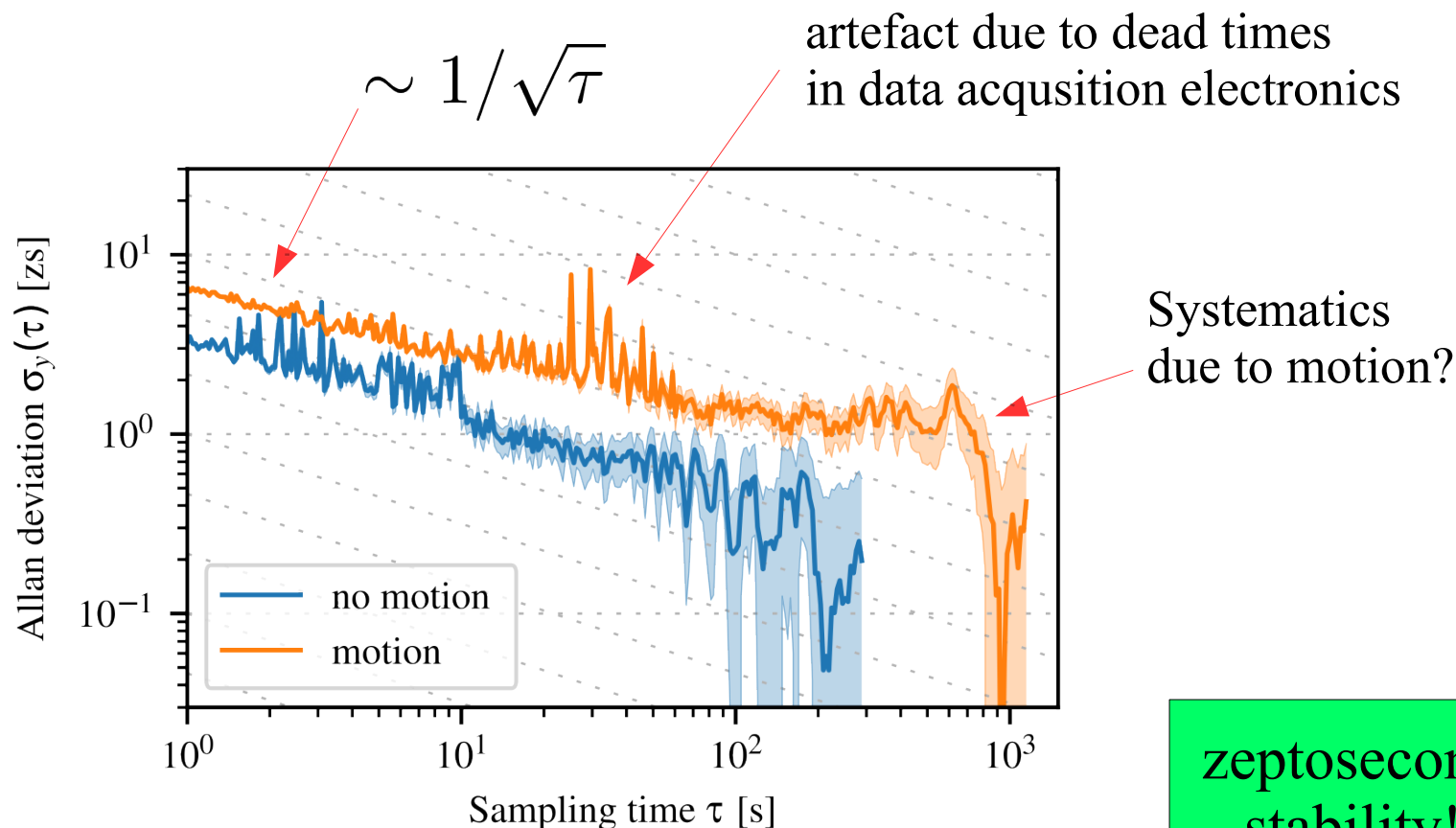
- Characteristic alternating intensities observed
- Fast oscillations due to additional beating with two resonances

How stable is the phase control?

Allan deviation: Measure for stability

- ▶ Split data into N equal parts of duration τ
- ▶ y_i is the phase measured in interval i
- ▶ RMS instability of two measurements τ apart

$$\sigma_y(\tau) = \left(\frac{1}{2(N-1)} \sum_{i=1}^{N-1} (y_{i+1} - y_i)^2 \right)^{\frac{1}{2}}$$



Systematics due to motion?

zeptosecond stability!

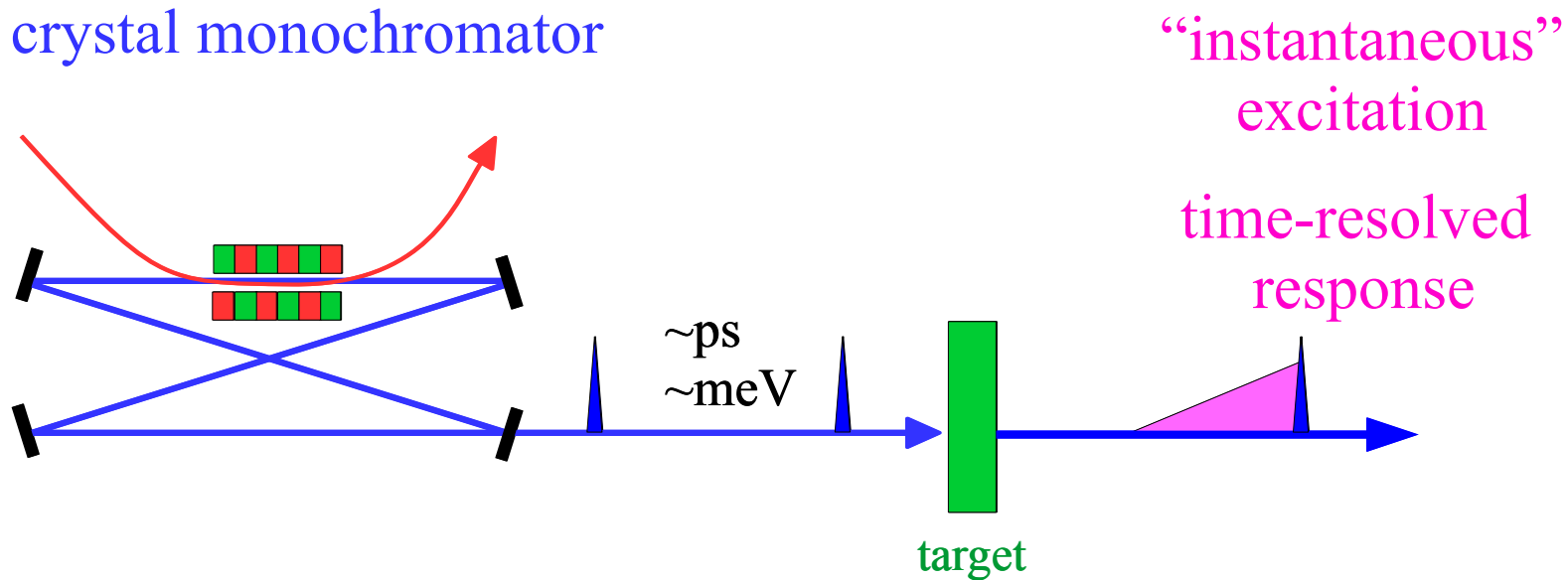
X-ray period
287 zs

Measurement
time ~ 45 min

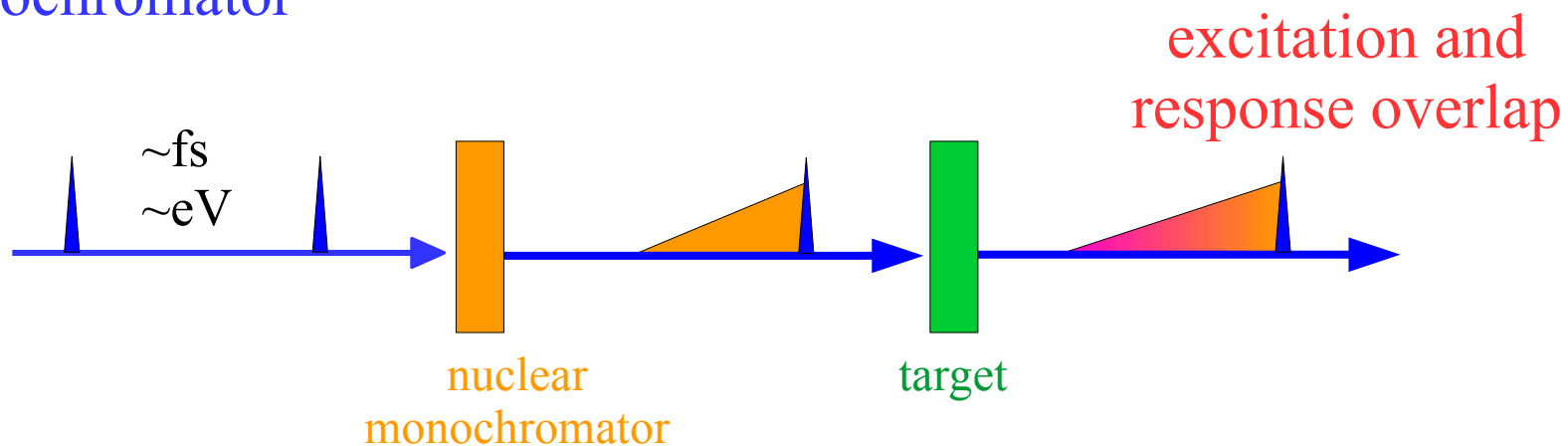


Temporal pulse structure

XFEL / crystal monochromator



Nuclear monochromator



Can we enter the non-linear regime?

Synchrotron: $\frac{0.01 \text{ Photons @ } 14.4\text{keV}}{100\text{ps bunch} \times (\mu\text{m})^2 \times \Gamma} \Rightarrow I \sim 10^2 \frac{\text{W}}{\text{cm}^2}$

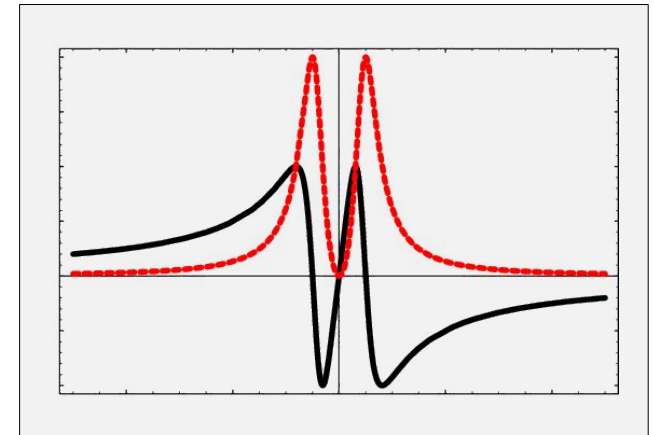
Seeded XFEL: $\frac{10^3 \text{ Photons @ } 14.4\text{keV}}{10\text{fs bunch} \times (\mu\text{m})^2 \times \Gamma} \Rightarrow I \sim 10^{10} \frac{\text{W}}{\text{cm}^2}$

EIT case: Kerr effect

$$n = n_0 + I_P n_2 \quad \chi = \chi^{(1)} + 3I_P \chi^{(3)}$$

$$\chi^{(3)} = 4.3 \times 10^{-22} \text{m}^2/\text{V}^2$$

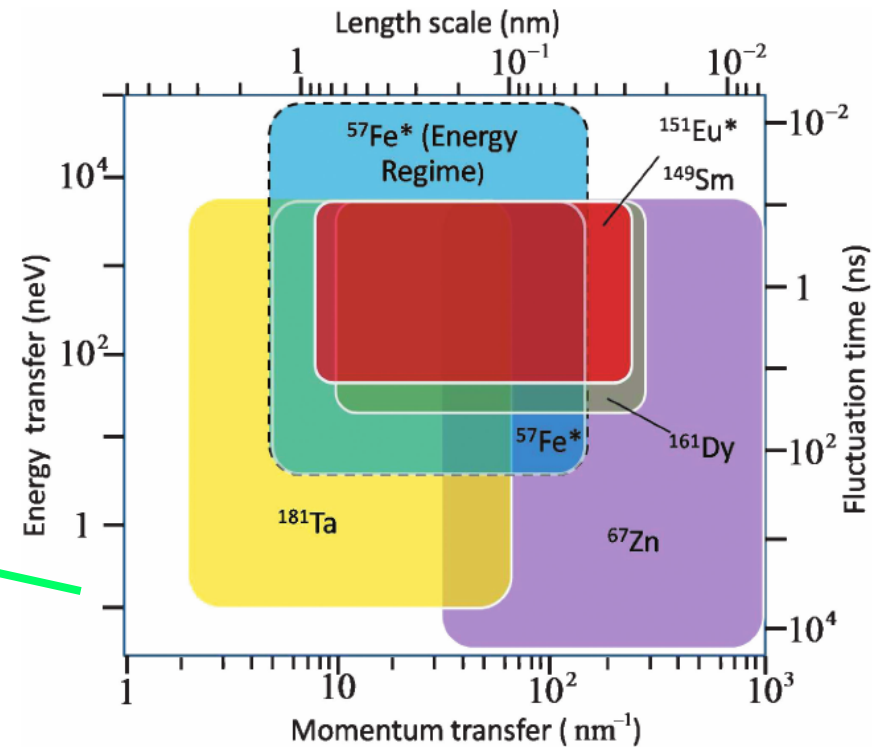
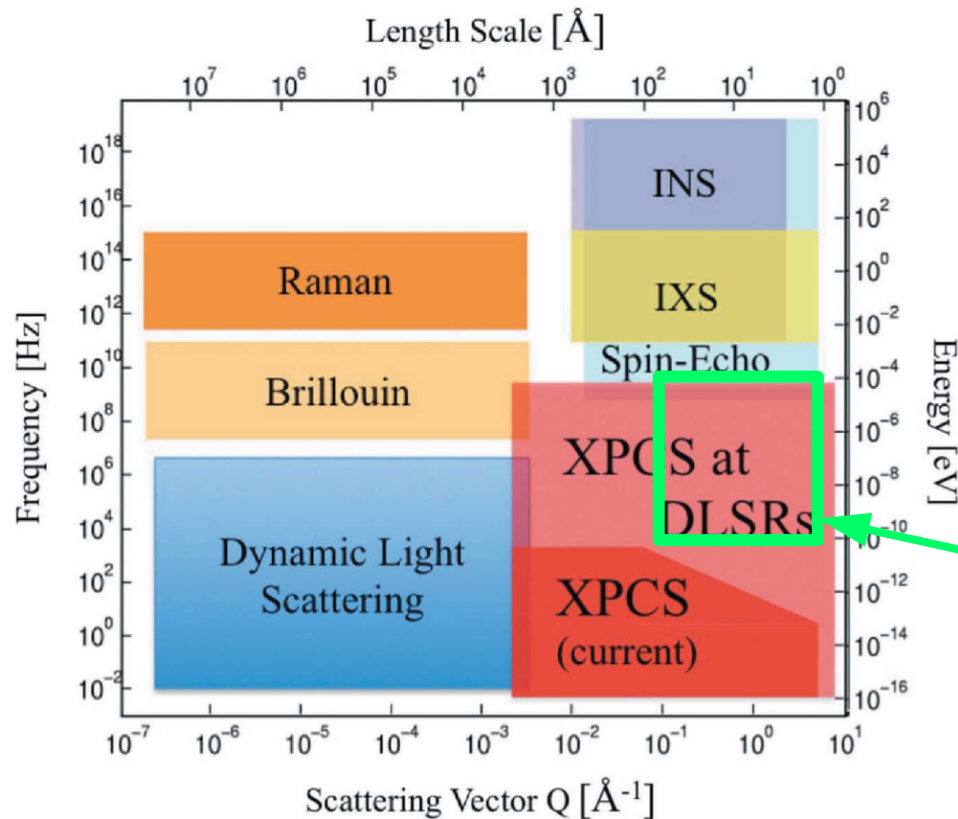
$$\Rightarrow n_2 I_P \approx 10^{-7} \text{ for } 10^8 \text{W}/\text{cm}^2$$



nonlinear phase shift \sim linear index achievable with seeded FEL
EIT: no linear absorption, strong enhancement via advanced schemes possible

The temporal gap: nanoseconds to milliseconds

Time domain interferometry with nuclear resonances



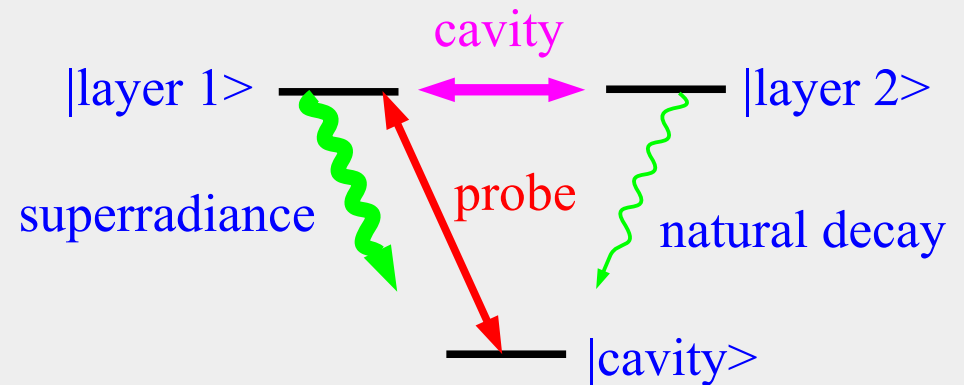
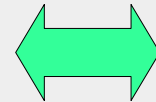
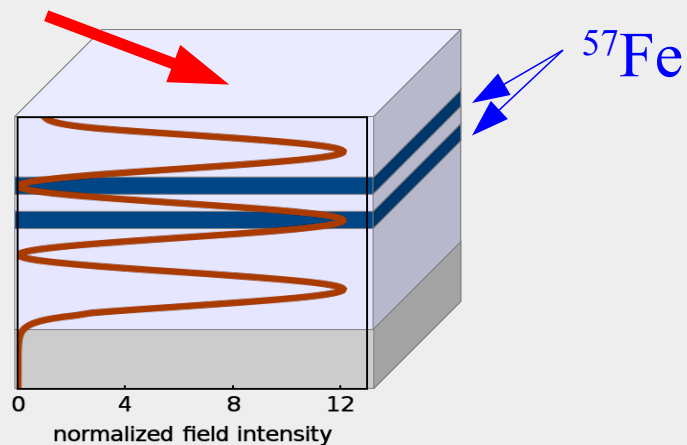
PETRA IV at DESY,
J. Synchrotron Rad. 25, 1277 (2018)

Saito et al., Appl. Phys.
Express 2, 026502 (2009)

+ TDI proposed for applications at XFELs in SwissFEL science case



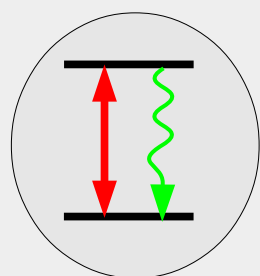
Design of artificial quantum systems



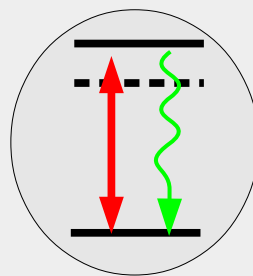
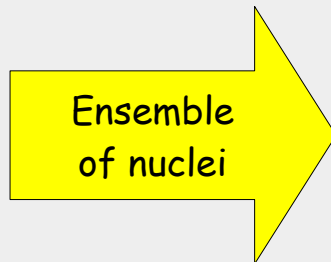
- create 3-level system otherwise not available
- replace external control field by cavity
- reduce decoherence via superradiance timescale

Röhlsberger et al, Nature 2012

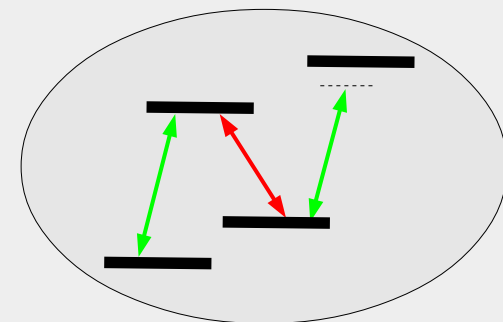
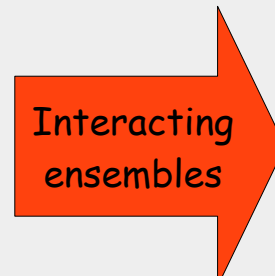
Next step: systematic approach to design level schemes



Single nucleus



Tunable two-level system



Tunable multi-level system

First steps: Longo, Keitel, Evers,
Sci. Rep. 6, 23628 (2016)

→ “reverse engineering” of superradiance
to determine required ensemble ensemble properties